

$$\begin{aligned} & \text{Гипотеза об отсутствии} \quad K = n - 1 \\ & V(x) = \frac{\omega^2 x^2}{2} \quad \sqrt{V} \\ & \langle x_0 | e^{-H(\tau_f - \tau_i)} | x_i \rangle = N \int_{-\infty}^{\infty} dx(x) e^{-\frac{\omega^2 x^2}{2}} \quad \xrightarrow{\text{так как}} \\ & \frac{dx}{dt} = \frac{d}{dt} \left(\frac{x - x_i}{\omega} \right) = \frac{dx(x)}{\omega} = \frac{dx}{dt} \quad \xrightarrow{\text{так как}} \\ & \langle x_0 | e^{-H(\tau_f - \tau_i)} | x_i \rangle = \sum_{k=1}^n e^{-E_k \tau_f} \quad \langle x_i | E_k \rangle \langle E_k | x_i \rangle \\ & \rightarrow -\omega = -\frac{d}{dt} \quad \xrightarrow{\text{так как}} e^{-E_k \tau_f} \langle x_i | 1 \rangle \sim e^{-E_k \tau_f} \quad \dots \\ & S_E = + \infty \quad S_E = \int dx \left(\left(\frac{dx}{dt} \right)^2 + \frac{\omega^2 x^2}{2} \right). \end{aligned}$$

$$\begin{aligned} \int d\tilde{x} e^{-\frac{f(\tilde{x})}{\delta^2}} \Big|_{g \rightarrow 0} &= \int d\tilde{x} e^{-\frac{f(\tilde{x}_0)}{\delta^2}} - \frac{f''(\tilde{x}_0)}{2\delta^2} (\tilde{x} - \tilde{x}_0)^2 = \\ f'(\tilde{x}_0) &= 0 \quad \leftarrow \text{согласно условию} \quad \left(\begin{array}{c} \dots \\ \dots \\ \dots \\ \dots \\ \dots \end{array} \right) \\ &= \frac{\sqrt{2\pi g}}{\sqrt{f''(\tilde{x}_0)}} e^{-\frac{f(\tilde{x}_0)}{\delta^2}} \\ \int d^n \tilde{x} e^{-\frac{f(\tilde{x})}{\delta^2}} &= \int d^n \tilde{x} e^{-\frac{f(\tilde{x}_0)}{\delta^2}} - \frac{\delta \tilde{x}_0}{\sqrt{2\pi g}} \frac{\partial^2 f}{\partial \tilde{x}^2} \frac{\delta \tilde{x}_0}{\delta^2} = \\ \frac{\partial f}{\partial \tilde{x}} \Big|_{\tilde{x}_0} &= 0 \quad \leftarrow \text{согласно} \quad \tilde{x} = \tilde{x}_0 \quad \left(\begin{array}{c} \dots \\ \dots \\ \dots \\ \dots \\ \dots \end{array} \right) \\ &= \frac{\sqrt{2\pi g}}{\sqrt{a_1}} \dots \frac{\sqrt{2\pi g}}{\sqrt{a_n}} e^{-\frac{f(\tilde{x}_0)}{\delta^2}} \\ &= \frac{(2\pi g)^{n/2} g^n}{\sqrt{\det(f''(\tilde{x}_0))}} e^{-\frac{f(\tilde{x}_0)}{\delta^2}} \end{aligned}$$

$$\begin{aligned} \langle x_0 = 0 | e^{-H(\tau_f - \tau_i)} | x_i = 0 \rangle &= \int_{\substack{\delta x \perp \tau_i - \delta x(\tau_f - \tau_i) \\ \frac{d^2 x}{dt^2} - \omega^2 x = 0 \\ x(\tau_0) = x_i(\tau_0) = 0}} \int dx \left(\frac{1}{2} \delta x \left(-\frac{d^2 x}{dt^2} + \omega^2 x \right) \delta x \right) \\ &= \int \frac{N}{\sqrt{\det(-\frac{d^2}{dt^2} + \omega^2)}} \left(\begin{array}{l} x_0(\tau_0) = x_i(\tau_0) = 0 \\ \int \delta x \cdot x_0(t) = \lambda x_0(\tau_0) \\ \text{дат } L = \prod L_i \end{array} \right) \end{aligned}$$

$$\begin{aligned} \text{д) Теорема Гамильтона - Якоби} \quad \dot{x}^2 = -\frac{d^2}{dt^2} + V(x) \\ \left[\begin{array}{l} \frac{d}{dt} \Psi(t) = 0 \\ \Psi(\tau_i) = 1 \end{array} \right] \Rightarrow \dot{x} + \dot{L} = \frac{N}{\lambda} \frac{d}{dt} \Psi(t) \quad \text{или } \Psi \text{ не зависит от } t \\ \text{Очевидно:} \quad \left\{ \begin{array}{l} \left(-\frac{d^2}{dt^2} + \omega^2 \right) \Psi = 0 \\ \frac{d\Psi}{dt}(\tau_i) = 1 \end{array} \right\} \quad \dot{x} = \frac{1}{\lambda} \sin(\omega(t - \tau_i)) \\ \det(-\frac{d^2}{dt^2} + \omega^2) = \frac{N}{\lambda} \sin(\omega(\tau_f - \tau_i)) \\ \langle x=0 | e^{-H(\tau_f - \tau_i)} | x=0 \rangle = \frac{N \lambda}{\sqrt{\sin(\omega(\tau_f - \tau_i))}} \quad \text{или } \lambda = \frac{N \lambda}{\sqrt{\tau_f - \tau_i}} \\ \langle x_0 | e^{-H(\tau_f - \tau_i)} | x_i \rangle = \frac{C - \frac{(x_f - x_i)^2}{2(\tau_f - \tau_i)}}{\sqrt{2\pi(\tau_f - \tau_i)}} \quad \text{или } C = \frac{1}{\sqrt{\tau_f - \tau_i}} \\ \langle x_{\infty} | e^{-H(\tau_f - \tau_i)} | x_{\infty} \rangle = \sqrt{\frac{\omega}{2\pi}} \frac{1}{\sqrt{\sin(\omega(\tau_f - \tau_i))}} \xrightarrow{\tau_f \rightarrow \infty} |\langle 0 | x=0 \rangle|^2 e^{-\frac{\omega}{2} E_0 \tau_f} \end{aligned}$$

$$\begin{aligned} V &= \frac{g^2}{2} \left(x^2 - \frac{a^2}{g^2} \right)^2 \\ &= \frac{1}{2g^2} \left((xg)^2 - a^2 \right)^2 = \frac{\nabla^2(xg)}{2g^2} \\ S_E &= \int dx \left(\frac{x^2}{2} + \frac{V(x)}{2} \right) = \frac{1}{2g^2} \int dx \left(\frac{g^2(xg)^2}{2} + \nabla^2(xg) \right). \end{aligned}$$

$$\langle x \rightarrow a | e^{-H(\tau_f - \tau_i)} | x = -a \rangle = N \int d\tilde{x}(x) e^{-\frac{S_E[\tilde{x}]}{2}}$$

согласно $\frac{d}{dt} \tilde{x}(x) = \nabla^2(\tilde{x}) = -(-\nabla)^2$

$$\begin{aligned} \tilde{x}_0 &= -a \\ &+ a \\ \tilde{x}_0 &= a + \ln\left(\frac{a}{\sqrt{2}}\right) \end{aligned}$$

$$\begin{aligned}
 & \omega = m_g \\
 & e^{-\int F(t^*) e^{-i\omega t^*} dt^*} \\
 & F(t^*) = \left(\frac{1}{1 + q^2/m_g^2} \right)^2 \\
 & \int d\theta e^{iqz \cos \theta} = \sin \theta d\theta = \\
 & 2\pi e^{\int_0^\pi dq \left(\frac{q^2}{1 + q^2/m_g^2} \right)^2} \int_0^\pi d\theta e^{iqz \cos \theta} = \sin \theta d\theta = \\
 & t = \frac{-\cos \theta}{\sin \theta} \quad d\theta = \sin \theta d\theta \\
 & 2\pi e^{\int_0^\infty dq \frac{q^2}{(1 + q^2/m_g^2)^2}} = 2\pi e \cdot \frac{1}{2} \quad \text{Im} \int_{-\infty}^{\infty} \frac{q e^{iqz}}{(1 + q^2/m_g^2)^2} \\
 & = \frac{8\pi e}{2} e^{-m_g z} = \rho(0) e^{-z/\alpha} \\
 & \frac{1}{2} 4\pi e i \lim_{q \rightarrow imy} \frac{d}{dq} \left[(1 + q^2/m_g^2)^2 \right] \frac{e^{izq}}{(1 + q^2/\alpha^2)^2}
 \end{aligned}$$

$$\int_{-\infty}^{\infty} \frac{qe^{iqz}}{(1+\frac{q}{m^2})^2} = -\int_{m^2}^{\infty} \frac{qe^{iqz}}{(1+\frac{iq}{m})^2(1-\frac{iq}{m})^2} \Theta$$
$$\begin{aligned} & \Theta - 2\pi i \cdot \lim_{q \rightarrow im} \frac{d}{dq} \left(\frac{qe^{iqz}}{(1-\frac{iq}{m})^2} \right) m^2 = \\ &= -2\pi m^2 i \lim_{q \rightarrow im} \left(\frac{e^{iqz}}{(1-\frac{iq}{m})^2} + \frac{q i e^{iqz}}{(1-\frac{iq}{m})^2} + \frac{2q e^{iqz} \frac{i}{m}}{(1-\frac{iq}{m})^3} \right) = \\ &= -2\pi m^2 i \left(\frac{e^{imz}}{4} - \frac{m^2 e^{-imz}}{4} - \frac{e^{-imz}}{4} \right) = i \frac{\pi}{2} m^3 e^{-imz} \\ & \frac{2\pi m^2 e^{-imz}}{2} = f(\omega) e^{-imz} \end{aligned}$$

$$\begin{aligned} & \langle p' | \bar{\psi}_r(q) | p \rangle = \\ & = \bar{u}_{p'} \{ \gamma^\mu F_1(q^2) - \frac{F_2(q^2)}{2m} S^\mu_\nu q_\nu \} u_p \end{aligned}$$

$$= \bar{u}_p \times \gamma^m F_1(q^2) - \frac{\gamma_2(q)}{2m} S^m q_v \} u_p$$

$$= \overline{u}_P \left\{ \delta^r F_m(q^2) + 2M [F_e(q^2) - F_m(q^2)] \frac{P^r}{P^2} \right\}$$

$$\delta_{\mu\nu}(\delta_{\mu\rho} - \delta_{\nu\rho})\epsilon^\rho = 0$$

$$\begin{aligned} \partial_\mu \bar{\psi}^\dagger \gamma^\mu \bar{\psi} &= \frac{i}{2} [\partial_\mu \bar{\psi}^\dagger - \bar{\psi}^\dagger \partial_\mu] (\bar{\psi}^\dagger \gamma^\mu \bar{\psi}) = \frac{i}{2} \bar{\psi}_\mu \hat{P}' - \frac{i}{2} \bar{\psi}_\mu m \\ \bar{\psi}^\dagger = \bar{P}' - p_0 \sqrt{\frac{i}{2} (\partial_\mu \bar{P}' - m \partial_\mu + p_\mu)} &= -\frac{i}{2} (\bar{\psi}_\mu \gamma^\mu - \bar{\psi}_\mu \partial_\mu) (\bar{\psi}^\dagger \gamma^\mu \bar{\psi}) \\ &= -\frac{i}{2} \bar{\psi}_\mu \hat{P}' - \cancel{\frac{i}{2} \bar{\psi}_\mu m} - \frac{i}{2} (m - \hat{P}') \bar{\psi}_\mu \end{aligned}$$

$$\sum_i \hat{P} \delta_\mu = P^0 (2 g_{\mu 0} - \delta_\mu \delta_0) = \cancel{2k p_r - \cancel{\delta_\mu m}} \quad P^r = P^1 r + \dots$$

$$z^{m\mu} q^\nu = \frac{1}{2} \left(p_\mu + p_\nu - 2m \delta^{\mu\nu} \right) \delta_{\mu\nu} p^\nu = \frac{u}{2} (k_0^2)$$

$$\delta^r F_m + 2m \left[F_e - F_m \right]$$

$\partial_\mu F_i(q^2)$

$$\delta_F F_1(q^2) = m \frac{F_2(q^2)}{2m^2} \left(p_\mu' + p_\mu^- \right)$$

$$\delta_m(F_1 + F_2) = \frac{2}{\pi} \int_{-\infty}^{\infty} \frac{e^{-imx}}{x^2 + m^2} dx$$

$$-\frac{4mF_2}{\gamma(\vec{p}^2 + \alpha^2)} P^\mu + \partial_\mu(F_\nu +$$

$$\overline{u}_p \left\{ \gamma^{\mu} F_m(q^2) + 2\mu(F_e - F_m) \frac{p^\mu}{p^2} \right\} u_p (-) \frac{(p^r + p', t)}{(q^2 - 4p^2) - i\epsilon}$$

$$F_m = F_1 + F_2 \quad \frac{q^2}{4M^2}$$

$$\partial \sigma^m < p' | \mathcal{I}_n(\underline{\lambda}) | p \rangle = \underbrace{(c_1 + c_2 \sigma^m + c_3 \sigma^{m+1} \sigma^m)}_{\text{operator}} \langle \sigma^m | p' | \mathcal{I}_n(\underline{\lambda}) | p \rangle$$

$$\cdot f(x) = e^{iP_x}$$

$$\int \langle p' | J_\mu^{\text{em}}(x) | p \rangle dx e^{iQx}$$

$$= \int \langle p' | e^{i\hat{P}_x} J_\mu(x) e^{-i\hat{P}_x} | p \rangle dx e^{iP_x}$$

$$\int dx e^{iQx} e^{ip'x - ipx}$$

$$\langle p' | J_\mu(x=0) | p \rangle$$

$$J_\mu(q)$$

$$\langle p' | j_{\mu}^{\text{em}}(x) | p \rangle \quad F(0)=0$$

$$j_\mu(q)$$

$$F_R(0)=1$$

