

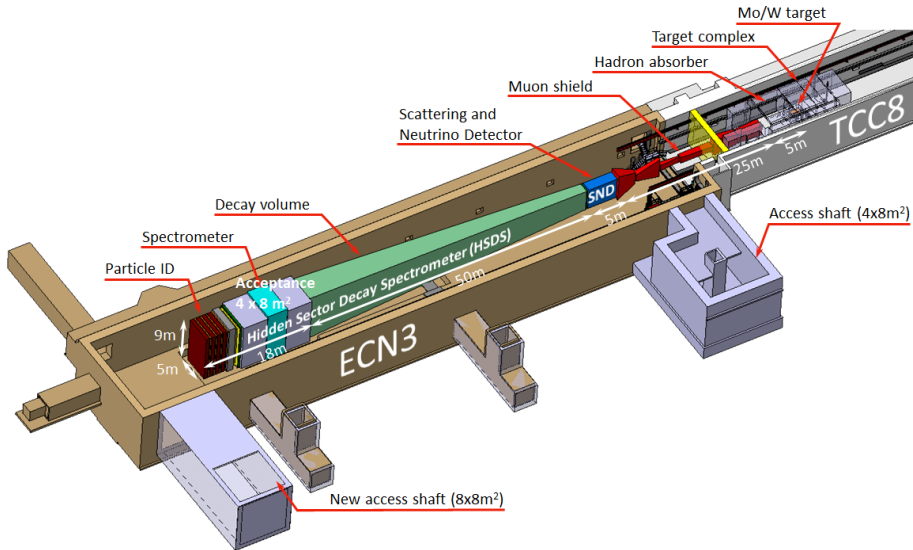
# FIMPs @ SHiP

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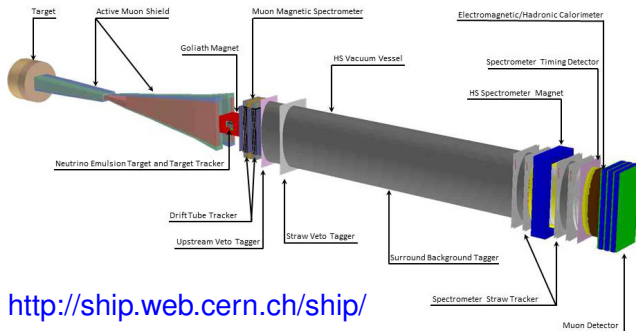
**Scientific Session devoted to 70th V.A.Rubakov Anniversary  
Nuclear Physics Division, Physics Department of RAS**

**RAS, Moscow, Russia**



# Towards a dedicated experiment

- $\nu$ MSM: T.Asaka, S.Blanchet, M.Shaposhnikov (2005), T.Asaka, M.Shaposhnikov (2005), D.G., M.Shaposhnikov (2007)
- direct tests of  $\nu$ MSM: D.G., M.Shaposhnikov (2007)
- proposal for direct searches, to European Strategy Group, 2012 D.G., M.Shaposhnikov
- sketch of realistic experiment S.Gninenko, D.G., M.Shaposhnikov (2013)
- Expression Of Interests: Proposal to Search for Heavy Neutral Leptons at the SPS W. Bonivento, ... D.G., et al, 1310.1762
- Technical Proposal and Physics Paper 1504.04956, 1504.04855
- included in the CERN GreyBook (2016) approved in 2024 !! 46 institutes from 16 countries



<http://ship.web.cern.ch/ship/>

# Standard Model + GR : Major Problems

Gauge and Higgs fields (interactions):  $\gamma, W^\pm, Z, g, G,$  and  $h$

Three generations of matter:  $L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}, e_R; Q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}, d_R, u_R$

- Describes all experiments dealing with
  - ▶ electroweak and strong interactions
- Does not describe (PHENO)
  - ▶ Neutrino oscillations
  - ▶ Dark matter ( $\Omega_{DM}$ )
  - ▶ Baryon asymmetry ( $\Omega_B$ )
  - ▶ Why the Universe is flat and homogeneous?
  - ▶ Where did the matter perturbations come from?
- (THEORY)
  - ▶ Dark energy ( $\Omega_\Lambda$ )
  - ▶ Strong CP-problem
  - ▶ Gauge hierarchy
  - ▶ Quantum gravity
  - ▶ Quantization of electric charge
  - ▶ Why 3 generations?
  - ▶ Why  $Y_e \ll Y_\mu \ll \dots \ll Y_t$

## Also, generally, we have anomalies...

- LEP (95 GeV in  $b\bar{b}$ ), CMS + ATLAS ( $\gamma\gamma$ ,  $\tau^+\tau^-$ )
- $B \rightarrow (K, K^*)\tau^+\tau^-$ , ...  $B \rightarrow K\pi\pi$ , etc
- muon  $g-2$  ... gone?
- LSND, MiniBooNE, reactor (gone?), Gallium, etc
- Superheavy BH in the present galaxies, in the early galaxies, BH binaries, etc
- GW-signal in Pulsar Timing Arrays...
- Extragalactic magnetic fields...

# Light NP: a logically possible option

- All the new particles are at (below)  $E_{EW}$   
then quantum contributions to  $m_h \sim E_{EW}$  are safe
  - Why so far no evidences for such light New Particles ?
  - They are only feebly coupled to the Standard Model
    - ▶ they are SM gauge singlets
    - ▶ new Yukawa-type couplings ?
    - ▶ portal-like couplings ?
- (not a GUT)

# Three Portals to the hidden World

Renormalizable interaction including SM field and new (hypothetical) fields singlets with respect to the SM gauge group

Attractive feature: couplings are insensitive to energy in c.m.f., hence low energy experiments (intensity frontier) are favorable

- Scalar portal: SM Higgs doublet  $H$  and hidden scalar  $X$  the simplest dark matter

$$\mathcal{L}_{\text{scalar portal}} = -\beta H^\dagger H X^\dagger X - \mu H^\dagger H X$$

- Spinor portal: SM lepton doublet  $L$ , Higgs conjugate field  $\tilde{H} = \varepsilon H^*$  and hidden fermion  $N$  sterile neutrino !!

$$\mathcal{L}_{\text{spinor portal}} = -y \bar{L} \tilde{H} N$$

- Vector portal: SM gauge field of  $U(1)_Y$  and gauge hidden field of abelian group  $U(1)'$  hidden photon

$$\mathcal{L}_{\text{vector portal}} = -\frac{\varepsilon}{2} B_{\mu\nu}^{U(1)_Y} B_{\mu\nu}^{U(1)'}$$

# Most natural example: coupling to inflaton

If  $V(\phi)$  dominates by chance

Chaotic inflation, A.Linde (1983), A.Linde (1984)

$$\ddot{\phi} - \Delta\phi/a^2 + 3H\dot{\phi} + V'(\phi) = 0$$

for power-law potential at  $\phi > M_{Pl}$

$$V \simeq \text{const}$$

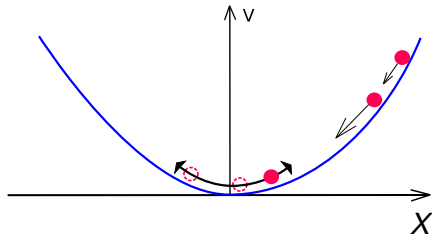
“slow roll” solution

$$H^2 = \frac{8\pi}{3M_P^2} V(\phi), \quad a(t) \propto e^{Ht}$$

valid while

slow roll conditions

$$M_P^2 \frac{V'''}{V} \ll 1, \quad M_P^2 \frac{V'^2}{V^2} \ll 1$$



$$-\beta H^\dagger H X^2 - \mu H^\dagger H X$$

Inflaton must couple  
to Standard Model fields

to reheat the Universe  
after inflation



# Renormalizable inflaton at GeV scale

0912.0390

$$S_{\text{XSM}} = \int \sqrt{-g} d^4x (\mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{ext}} + \mathcal{L}_{\text{grav}}),$$

$$\mathcal{L}_{\text{ext}} = \frac{1}{2} \partial_\mu X \partial^\mu X + \frac{1}{2} m_\chi^2 X^2 - \frac{\beta}{4} X^4 - \lambda \left( H^\dagger H - \frac{\alpha}{\lambda} X^2 \right)^2,$$

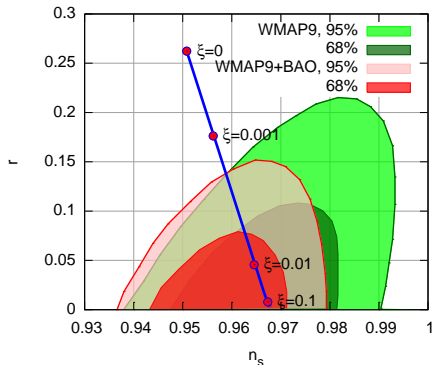
$$\mathcal{L}_{\text{grav}} = - \frac{M_{\text{P}}^2 + \xi X^2}{2} R,$$

inflaton mass

$$m_\chi = m_h \sqrt{\frac{\beta}{2\alpha}} = \sqrt{\frac{\beta}{\lambda \theta^2}}.$$

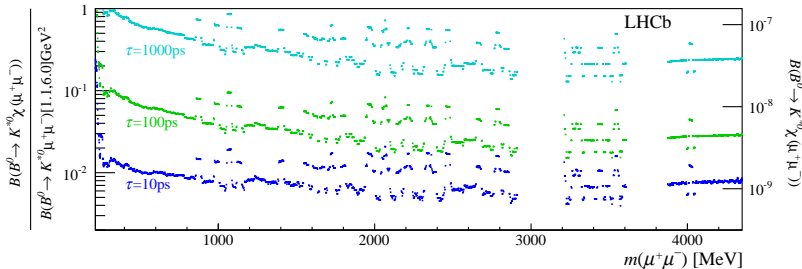
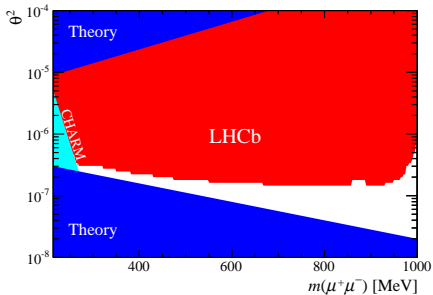
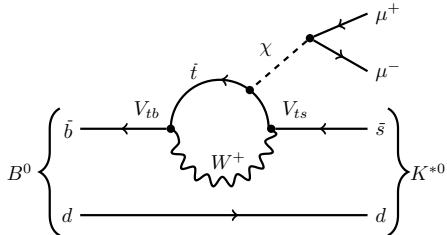
phenomenology is fixed by  
mixing with Higgs

$$\theta^2 = \frac{2\beta v^2}{m_\chi^2} = \frac{2\alpha}{\lambda}.$$



# Limits from LHCb

1508.04094



# Sterile neutrinos or Heavy Neutral Leptons

Most general renormalizable with 2(3...) right-handed neutrinos  $N_I$

$$\mathcal{L}_N = \bar{N}_I i \not{\partial} N_I - f_{\alpha I} \bar{L}_\alpha \tilde{H} N_I - \frac{M_{N_I}}{2} \bar{N}_I^c N_I + \text{h.c.}$$

## Parameters to be determined from experiments

9(7): active neutrino sector

2  $\Delta m_{ij}^2$ : oscillation experiments

3  $\theta_{ij}$ : oscillation experiments

1 CP-phase: oscillation experiments

2(1) Majorana phases:  $0\nu e e$ ,  $0\nu \mu \mu$

1(0)  $m_\nu$ :  ${}^3\text{H} \rightarrow {}^3\text{He} + e + \bar{\nu}_e$ , cosmology, ...

11:  $N = 2$  sterile neutrinos  
(works if  $m_\nu = 0$  !!!)

2: Majorana masses  $M_{N_I}$

9: New Yukawa couplings  $f_{\alpha I}$  which form

2: Dirac masses  $M^D = f \langle H \rangle$

3+1: mixing angles

2+1: CP-violating phases

4 new parameters in total

18:  $N = 3$  sterile neutrinos:

3: Majorana masses  $M_{N_I}$

15: New Yukawa couplings  $f_{\alpha I}$  which form

3: Dirac masses  $M^D = f \langle H \rangle$

3+3: mixing angles

3+3: CP-violating phases

9 new parameters in total

Profit: can suggest why neutrinos are so light,  $m_\nu \sim 0.1 - 0.01$  eV

# Sterile neutrino: a vast region of mass

Within the seesaw paradigm, as far as

$$m_a \sim \frac{f^2 v^2}{M_N^2} M_N \sim \theta^2 M_N$$

Any set

(mass scale  $M_N$ , Yukawa coupling  $f$ )

is viable

And with special tuning or symmetry larger (but not smaller) mixing

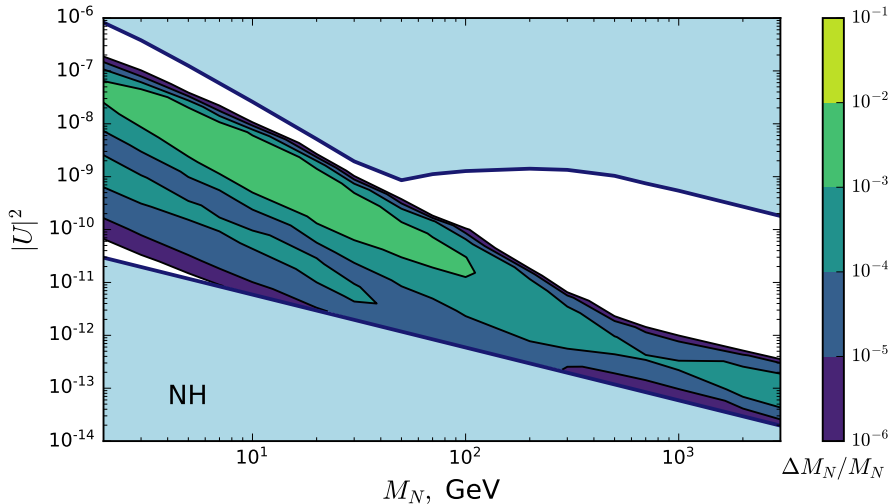
3 sterile neutrinos

is viable

$$\hat{m}_a \sim \hat{f}^T \frac{1}{\hat{M}_N} \hat{f} v^2$$

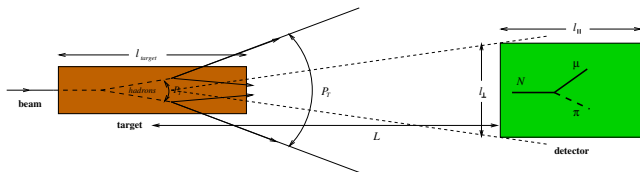
# Progress in Leptogenesis (no update on DM)

2008.13771



# Fixed target and similar

However for the feebly coupled light particle best place to show up is  
the intensity frontier fixed target experiment



## Variations and specifics

- dedicated (e.g. NA64) or working as by-product (e.g. T2K, DUNE)
- thin target (e.g. T2K, DUNE) or dump (e.g. NA64)
- decays or hits as the signature
- production by cosmic rays
- ...

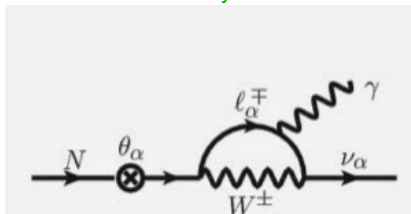
# Searches for sterile neutrino dark matter

$$m_a \sim \frac{f^2 v^2}{M_N^2} M_N \sim \theta^2 M_N$$

- **unstable**, but exceeding the age of the Universe if

$$\frac{\theta^2}{3 \times 10^{-3}} < \left( \frac{10 \text{ keV}}{M_N} \right)^5$$

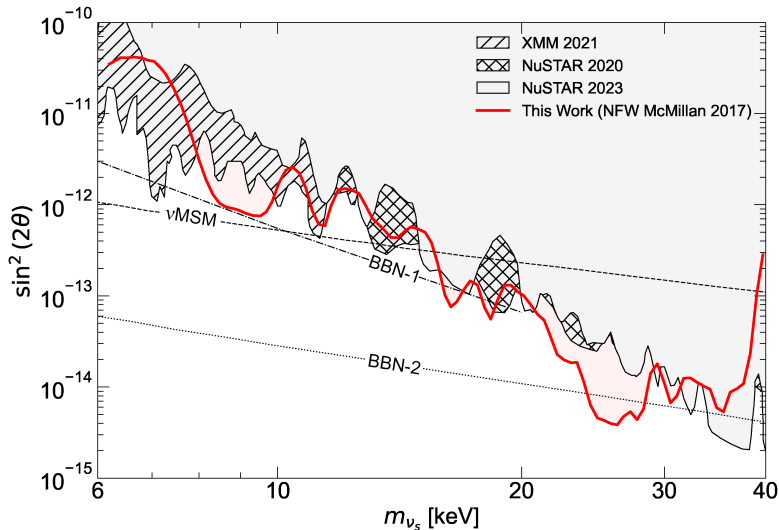
- **DM sterile neutrinos can be searched at X-ray telescopes because of two-body radiative decay**  
give limits in absence of the feature



a narrow line ( $\delta E_\gamma / E_\gamma \sim v \sim 10^{-3}$ )  
at photon frequency  $E_\gamma = M_N / 2$

$$\frac{\theta^2}{10^{-11}} \lesssim \left( \frac{10 \text{ keV}}{M_N} \right)^4$$

if redistributed equally... two different codes: `sterile-dm` and `dmpheno`



2405.17861



# Massive vectors (paraphotons)

NA64

Vector portal to a secluded sector:

one more  $U(1)'$  gauge group [spontaneously broken] in secluded sector

e.g. with Dark matter  $\Psi$

0711.4866

$$\mathcal{L}_{\text{DM+mediator}} = \bar{\Psi} \left( i\gamma^\mu \partial_\mu - e' \gamma^\mu A'_\mu - m_\Psi \right) \Psi - \frac{1}{4} A'_{\mu\nu} A'^{\mu\nu} + \frac{m_\gamma^2}{2} A'_\mu A'^\mu + \varepsilon A'_\mu \partial_\nu B^{\mu\nu}$$

when  $m_\Psi > m_\gamma \sim 1 \text{ GeV}$

- limit from BBN:

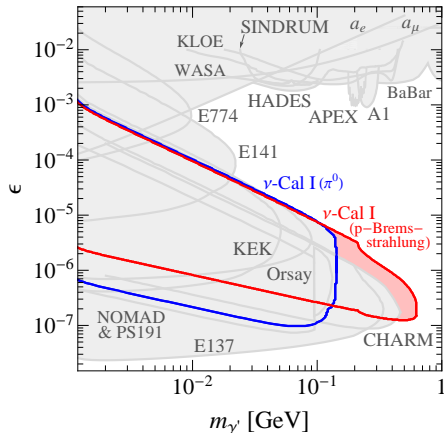
$$\tau_V < 1 \text{ s}, \implies \varepsilon^2 \left( \frac{m_\gamma}{1 \text{ GeV}} \right) \gtrsim 10^{-21}$$

- light for  $(g-2)$
- light for Pamela, Fermi, etc

Production by virtual photon  
Decay through virtual photon,  
 $V \rightarrow e^+ e^-, \mu^+ \mu^-, \text{ etc}$

$$\sigma \propto \varepsilon^2$$

$$\Gamma \propto \varepsilon^2$$



1311.5104

# Huge community: theory meets experiment

2305.01715

New models, reinterpretation of old data, new experimental constraints

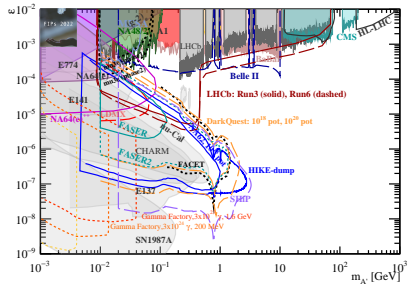
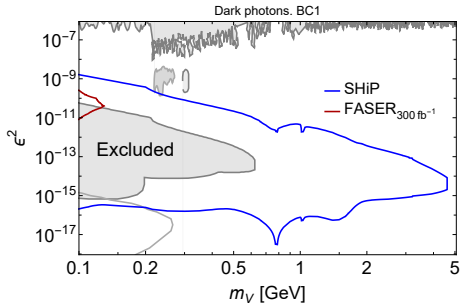
## Benchmark Models (BC#)

- 1 paraphoton (dark photon)
- 2  $B - L$  vector
- 3 millicharged
- 4 scalar ( $H \rightarrow SS$  at 0.01)
- 5 scalar (no  $H \rightarrow SS$ )
- 6 HNL mixed with  $\nu_e$
- 7 HNL mixed with  $\nu_\mu$
- 8 HNL mixed with  $\nu_\tau$
- 9 ALP coupled to photons
- 10 ALP coupled to fermions
- 11 ALP coupled to gluons

## Experiments and Proposals

- 1 T2K, T2HK, DUNE
- 2 NICA, FLAP
- 3 SHiP
- 4 DarkQuest, NA64, NA62-damp, PADME
- 5 FASER, FASER-2, SND
- 6 CODEX-b
- 7 FORMOSA, FLaRE
- 8 MilliQan
- 9 ANUBIS, MATHUSLA
- 10 SHADOWS, HIKE-K+, HIKE-dump
- 11 LDMX, M<sup>3</sup>, SBND

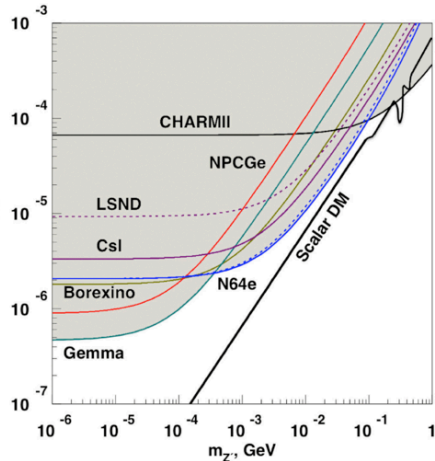
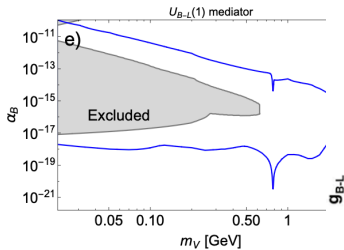
# BC1: paraphoton (dark photon) $V_\mu$



$$\Delta L_{int} = -\frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu} + \frac{m_V^2}{2} V_\mu V^\mu$$

decay

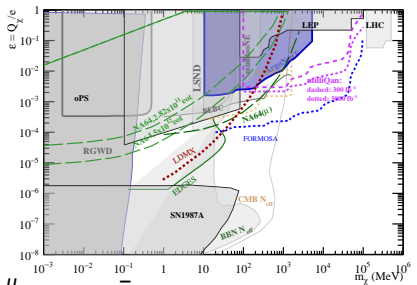
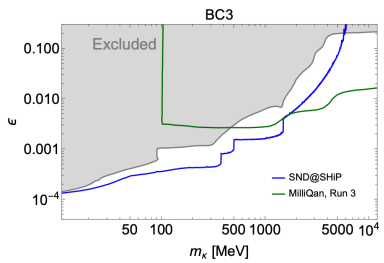
# BC2: vector $V_\mu$ of $U(1)_{B-L}$ gauge symmetry



$$\Delta L_{int} = -e_B V_\mu (j_B^\mu - j_L^\mu) + \frac{m_V^2}{2} V_\mu V^\mu$$

decay

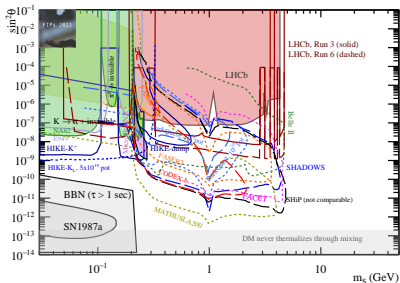
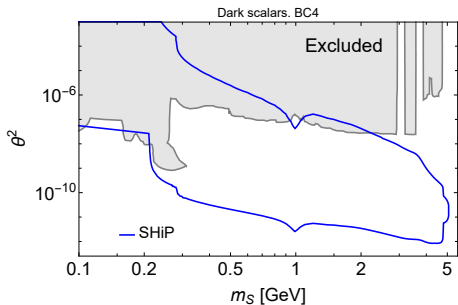
# BC3: millicharged particles, $q = \epsilon e$



$$\Delta L_{int} = -\epsilon e A_\mu \bar{\psi} \gamma^\mu \psi - m_\kappa \bar{\psi} \psi$$

hit

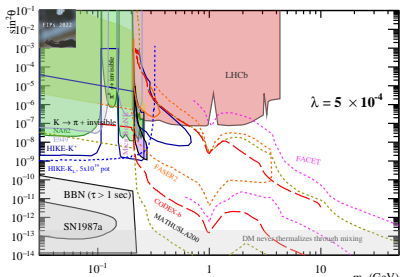
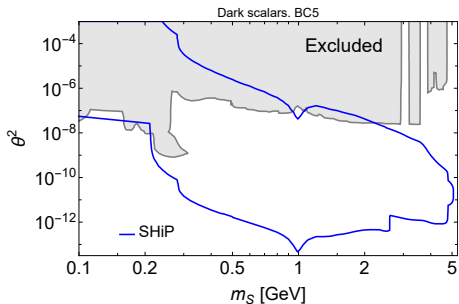
# BC4: scalar portal: only **mixing** with Higgs



$$\Delta L_{int} = \mu H^\dagger HS - \frac{m_S^2}{2} S^2 \rightarrow \theta hS - \frac{m_S^2}{2} S^2$$

decay

# BC5: scalar portal: cubic and quartic terms



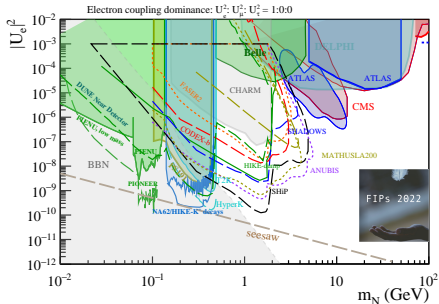
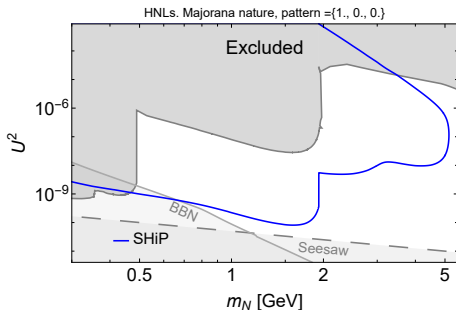
$$\Delta L_{int} = \mu H^\dagger H S - \frac{\lambda}{2} H^\dagger H S^2 - \frac{m^2}{2} S^2 \rightarrow \theta h S - \mu' h S^2 - \frac{m_S^2}{2} S^2$$

fixing  $\text{Br}(h \rightarrow SS) = 0.01$

double production from  $B$  – meson via  $\mu'$

decay

# BC6: Heavy Neutral Lepton $N$ coupled to $\nu_e$

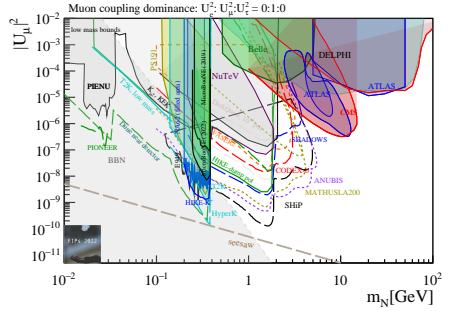
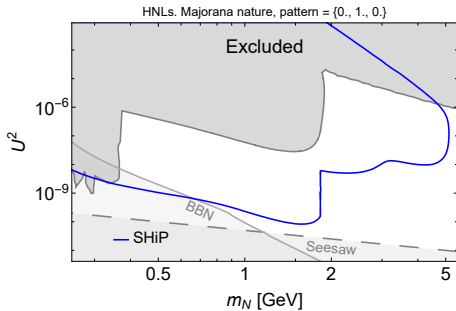


$$\Delta L_{int} = Y_{\alpha I} \bar{L}_\alpha \tilde{H} N_I : \nu_e \rightarrow \nu_e + U_e N$$

decay



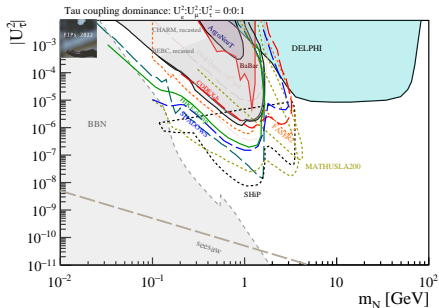
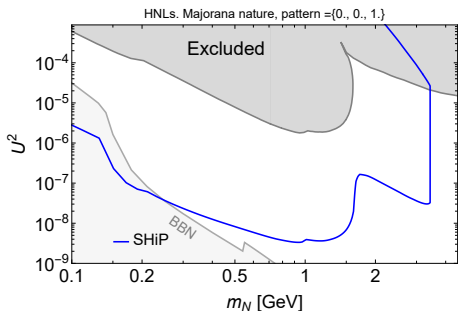
# BC7: Heavy Neutral Lepton $N$ coupled to $\nu_\mu$



$$\Delta L_{int} = Y_{\alpha l} \bar{L}_\alpha \tilde{H} N_l : \nu_\mu \rightarrow \nu_\mu + U_\mu N$$

decay

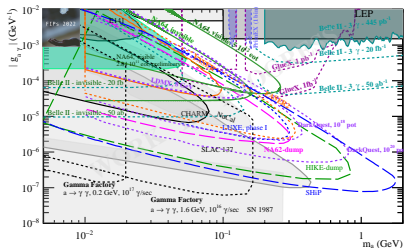
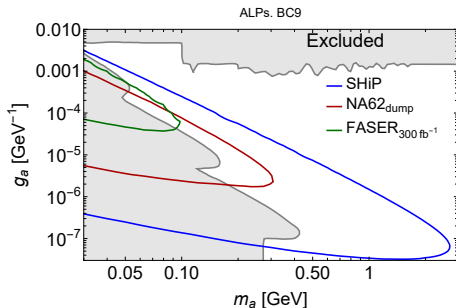
# BC8: Heavy Neutral Lepton $N$ coupled to $\nu_\tau$



$$\Delta L_{int} = Y_{\alpha I} \bar{L}_\alpha \tilde{H} N_I : \nu_\tau \rightarrow \nu_\tau + U_\tau N$$

decay

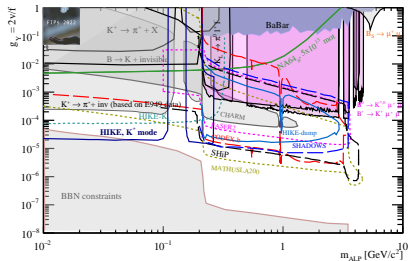
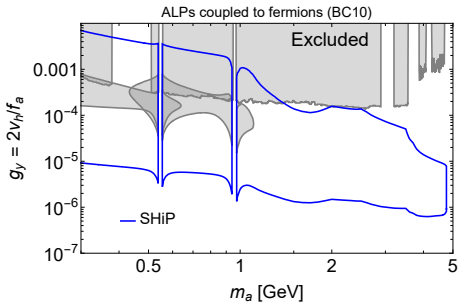
# BC9: Axion Like Particle $a$ coupled to photons



$$\Delta L_{int} = g_a a F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_a^2 a^2$$

decay

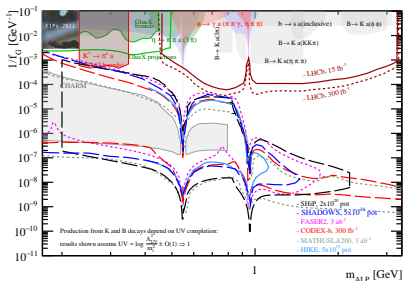
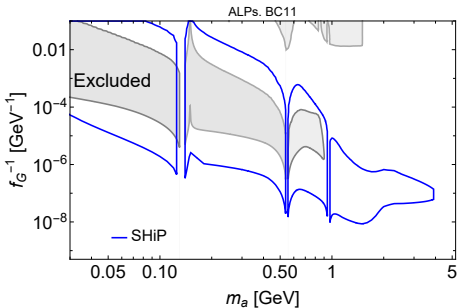
# BC10: Axion Like Particle $a$ coupled to fermions



$$\Delta L_{int} = g_y a \bar{\psi} \psi - \frac{1}{2} m_a^2 a^2$$

decay

# BC11: Axion Like Particle $a$ coupled to gluons



$$\Delta L_{int} = \frac{a}{f_g} G_{\mu\nu} G^{\mu\nu} - \frac{1}{2} m_a^2 a^2$$

decay

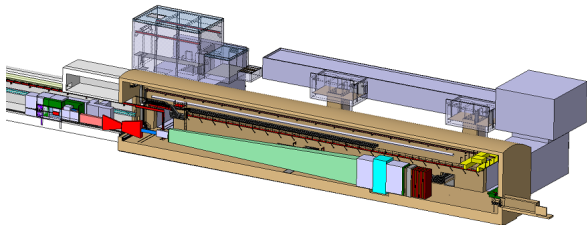
# SHiP: Search for Hidden Particles, $4 \cdot 10^{19}$ PoTs/year

Experiment at the SPS Beam Dump Facility to Search for Hidden Particles SHiP and the associated SPS Beam Dump Facility is a new general-purpose experiment in preparation at the SPS to search for "hidden" particles as predicted by a large number of models of Hidden Sectors that are capable of explaining for instance dark matter, neutrino oscillations, and the origin of the baryon asymmetry in the Universe. The experiment is design to search for any type of feebly interacting long-lived particles, among which are found e.g. heavy neutral leptons, dark photons, dark scalars, axion-like particles, and light supersymmetric particles - sgoldstinos, etc, as well as different types of Light Dark Matter. The high intensity of the SPS and in particular the large production of charm mesons and photons with the 400 GeV proton beam allow a comprehensive search at the MeV-GeV scale over many orders of magnitude in coupling. The detector incorporates two complementary apparatuses aimed at searching for hidden particles through both visible decays and through scattering signatures from recoil of electrons or nuclei. Moreover, the facility is ideally suited to study the interactions of tau neutrinos.



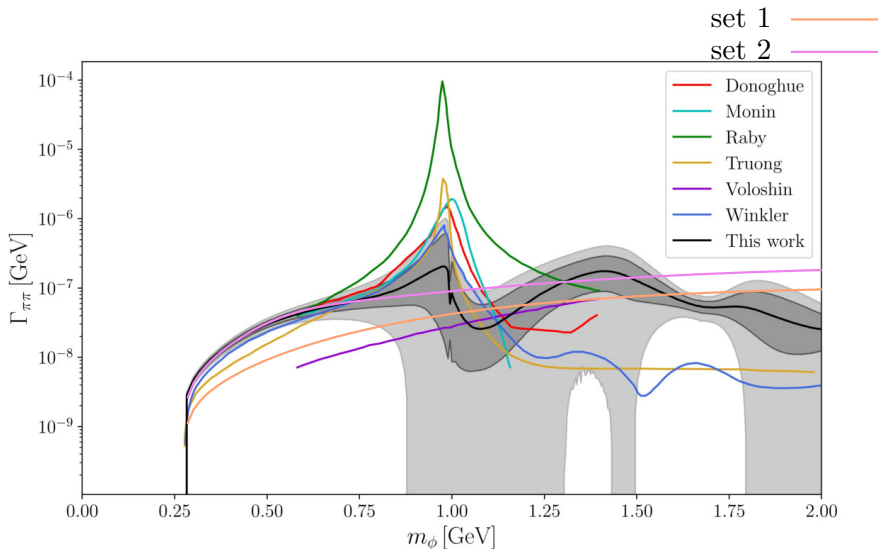
SHiP

Search for Hidden Particles



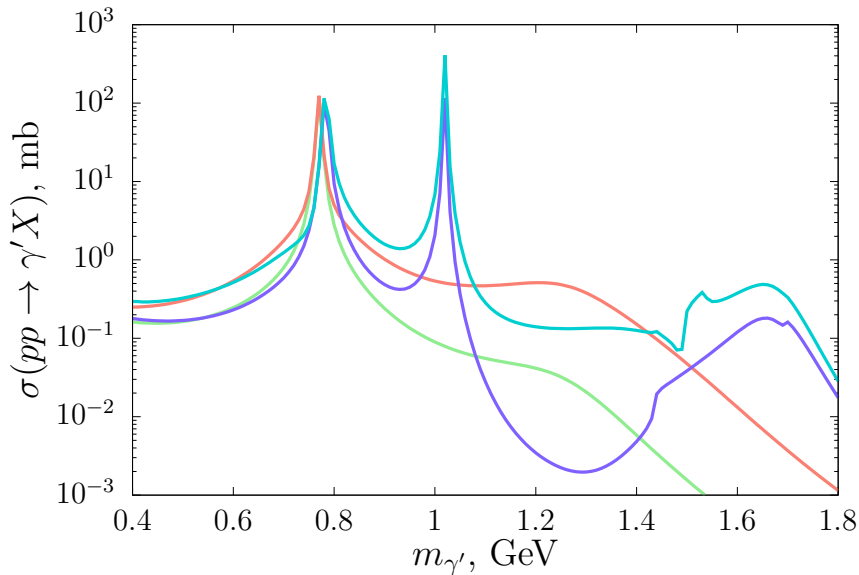


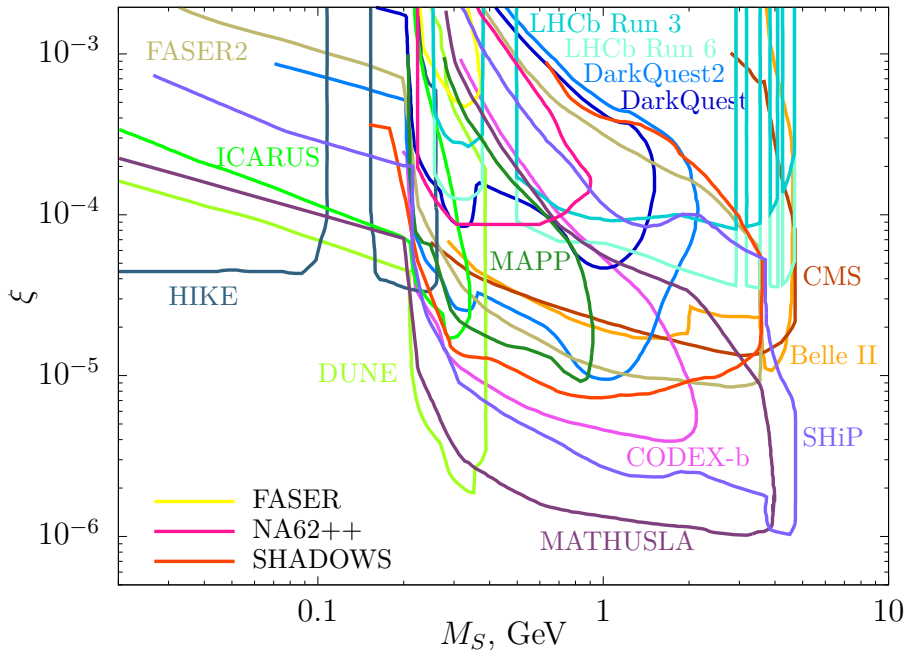
# Uncertainties 2407.13587 and comparison with GDA

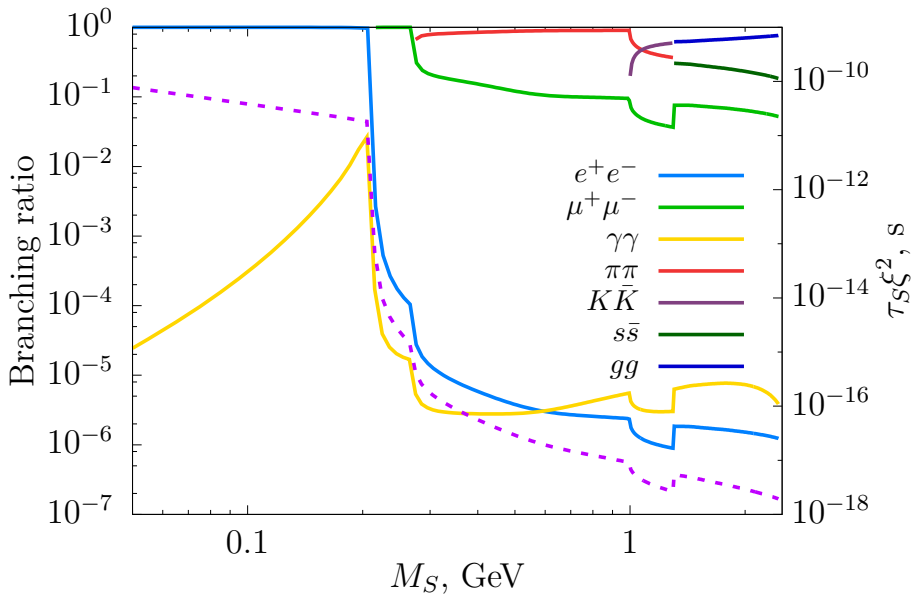




# different Dirac, Pauli form factors (0910.5589,2010.15872), talk by E. Kriukova

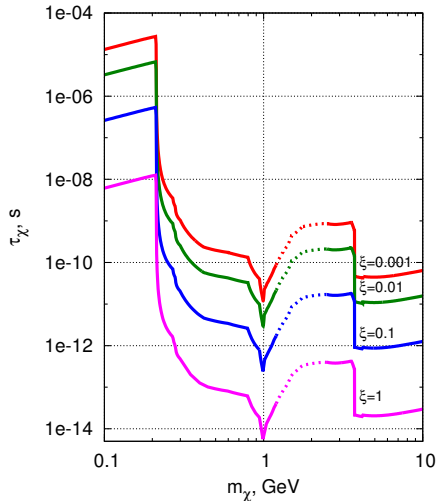
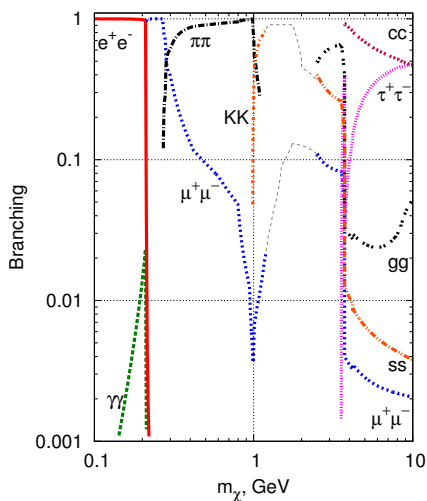






# QCD modes: claimed uncertainties upto $10^2$

1303.4395

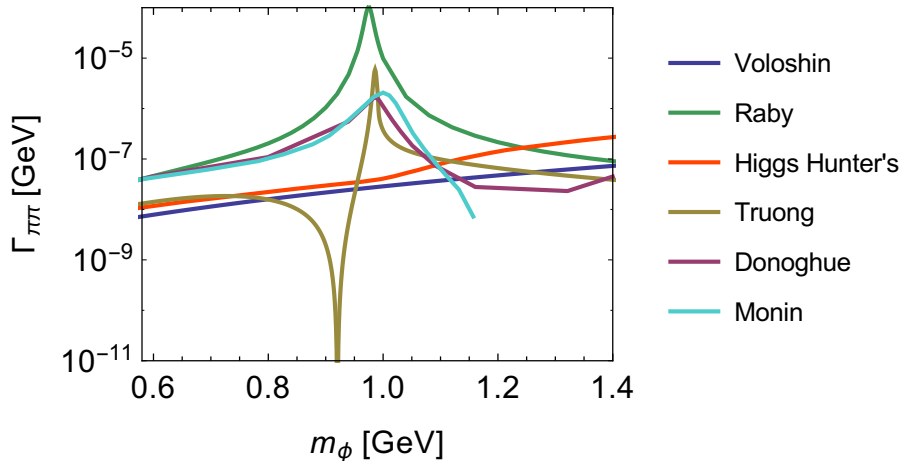


Interaction among the final hadronic states

following J.Donoghue, J.Gasser and H Leutwyler (1990)

# The estimates BSM people use

1809.01876



Physics: Final State Interaction:  $\pi\pi \rightarrow \pi\pi$

# sensitivity curves

2102.12143

Two different cases:

1  $I_{dec} = c\tau_X\gamma \ll L_{detector}$

$$N_{X \rightarrow chan}^{events} \propto Br_{chan}$$

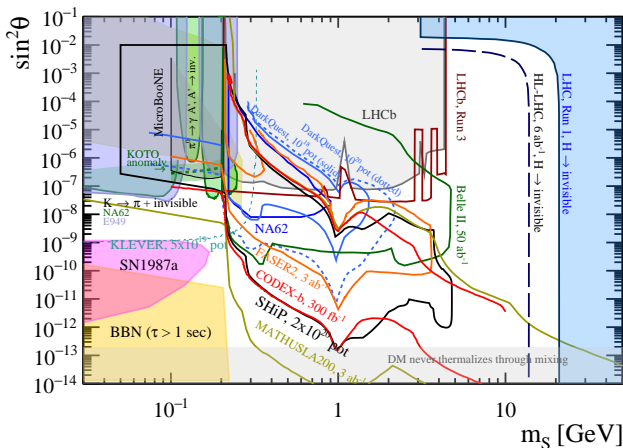
Smooth for pions, wiggly for muons

2  $I_{dec} = c\tau_X\gamma \gg L_{detector}$

$$N_{X \rightarrow chan}^{events} \propto \frac{Br_{chan}}{c\tau_X\gamma}$$

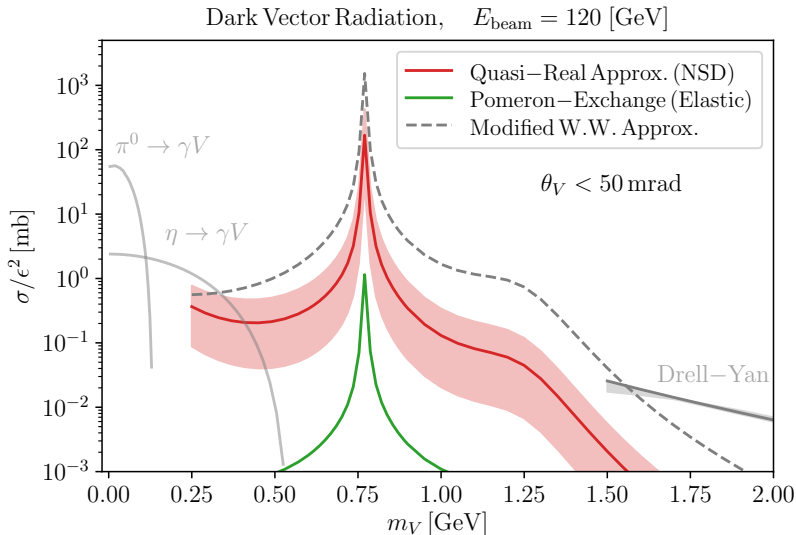
Wiggly for pions, smooth for muons

while pions dominate...

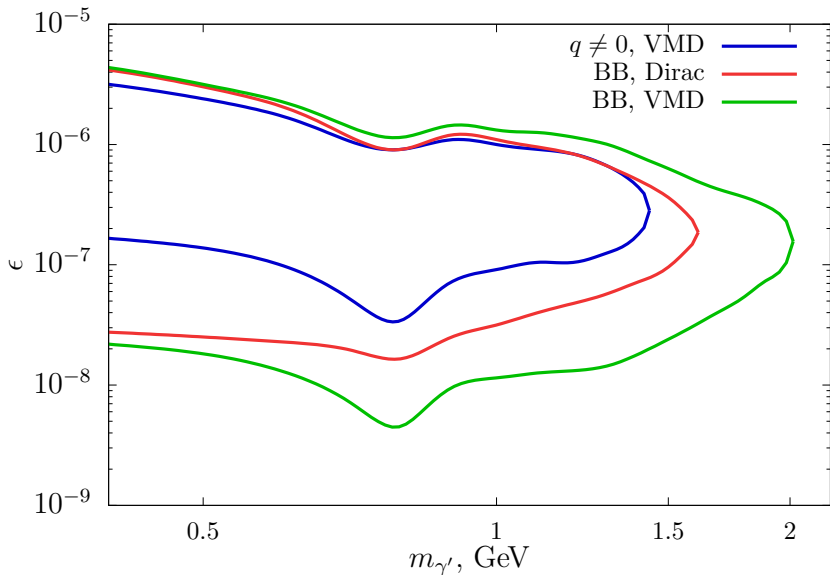


# Massive vectors: various attempts

2108.05900



# It changes the sensitivity... SHiP, 400 GeV





# Outcome for future experiments see talk by E. Kriukova

