

# Absorption of twisted photon by an electron in a strong magnetic field

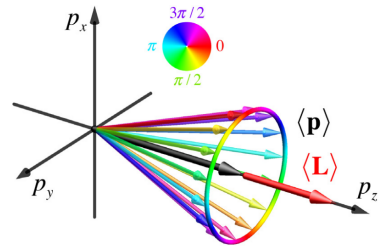
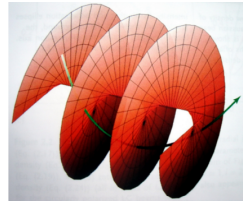
A.A Shchepkin, D.V. Grosman, I.I. Shkarupa and D.V. Karlovets

Faculty of Physics, ITMO University

# Twisted states

Twisted (vortex) states are **waves** with well-defined:

- longitudinal momentum,  $k_z$ ;
- absolute transverse momentum,  $\kappa = |\mathbf{k}_\perp|$ ;
- energy,  $E = \frac{\kappa^2 + k_z^2}{2m}$ ;
- angular momentum projection,  $l_z = m$ .

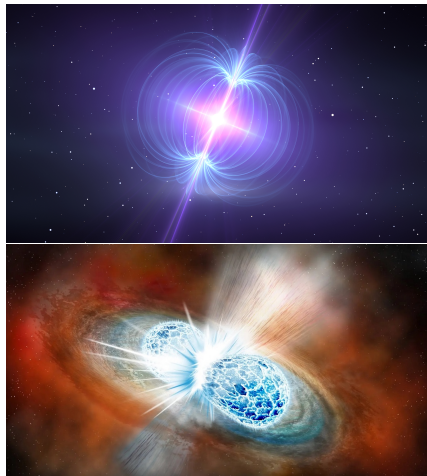


# Generation of vortices in strong magnetic field

$$(H_c = m_e^2 c^3 / e \hbar = 4.4 \times 10^{13} \text{ G})$$

Emission in strong magnetic field  
( $H \sim H_c$ ) is mostly twisted  
**(Maruyama+'22, Pavlov&Karlovets'24).**

Stellar nucleosynthesis. Radiative transfer in astrophysical environments such as neutron stars, kilonovae etc.



# Gauge of the vector potential

## Landau gauge

$$\mathcal{A}^\mu = \{0, 0, Hx, 0\}$$

$$\Phi(\mathbf{r}) \propto H_n(\tilde{x}) e^{-\tilde{x}^2/2} e^{i(p_z z + p_y y)}$$
$$\tilde{x} = (x - \rho_H^2 p_x) / \rho_H$$

Axial symmetry is not preserved.  
No orbital angular momentum.

## Symmetric gauge

$$\mathcal{A}^\mu = H/2 \times \{0, -y, x, 0\}$$

$$\Phi(\mathbf{r}) \propto \tilde{\rho}^l L_s^l(\tilde{\rho}) e^{-\tilde{\rho}^2/2} e^{ip_z z}$$
$$\tilde{\rho} = \rho / \rho_H$$

Axial symmetry is preserved.  
Orbital angular momentum ( $l$ )  
arises.

# Properties of Landau states in symmetric gauge

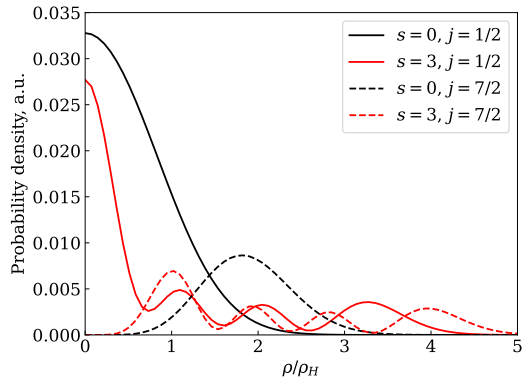
## R.m.s. radius

$$\langle \rho \rangle = \rho_H \sqrt{2s + j + 1} \quad (1)$$

## Energy

$$\varepsilon_{s,j}^2 = m_e^2 + p_z^2 + \frac{4}{\rho_H^2} \left( s + j + \frac{1}{2} \right), \quad (2)$$

$$s + j \geq -1/2$$

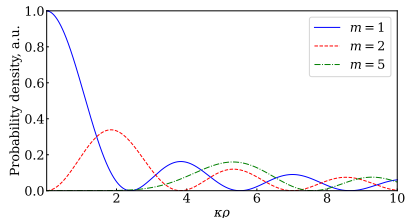
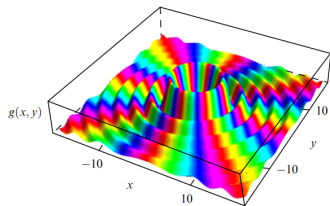


# Twisted photon

Emitted radiation by an electron in strong magnetic field is mostly twisted (**Maruyama'22**, **Pavlov&Karlovets'24**)

The photon is described with the Bessel shape (**Knyazev&Serbo'18**)

$$\psi_{\kappa m k_z} = J_m(\kappa\rho) \exp[i(m\varphi_r + k_z z)] \quad (3)$$



# Kinematics

## Conservation laws

$$\begin{aligned}\varepsilon_f &= \varepsilon_i + \omega, \\ p_{zf} &= p_{zi} + k_z, \\ j_f &= j_i + m.\end{aligned}\tag{4}$$

## Energy of the absorbed photon

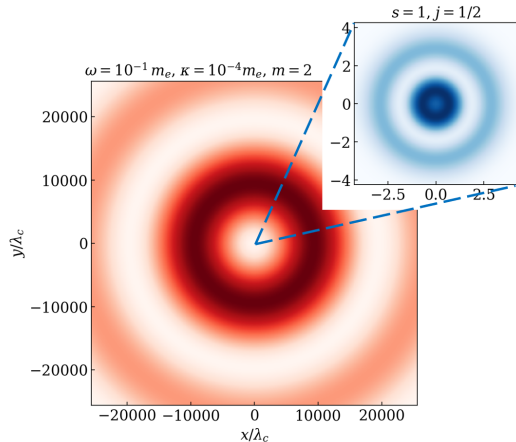
$$\omega = (s_f - s_i + m)\omega_c - \frac{\kappa^2}{2\varepsilon_{s_i, j_i}}, \quad \omega_c = \frac{eH}{\varepsilon_{s_i, j_i}}\tag{5}$$

# First order matrix element

$$S_{fi}^{(1)} = \langle e_f | \hat{S} | e_i, \gamma \rangle =$$

$$= -ie \int d^4x j_{fi}^\mu(x) A_\mu(x). \quad (6)$$

$$j_{fi}^\mu(x) = \bar{\Psi}_f(x) \gamma^\mu \Psi_i(x). \quad (7)$$





# Generalized cross section

## Cross section

$$d\sigma = \frac{dW}{\mathcal{L}} \quad (8)$$

## Absorption probability

$$dW = \left| S_{fi}^{(1)} \right|^2 \frac{L}{2\pi} dp_{zf} \Delta s_f \Delta j_f, \quad (9)$$

## Dependence on $\kappa$

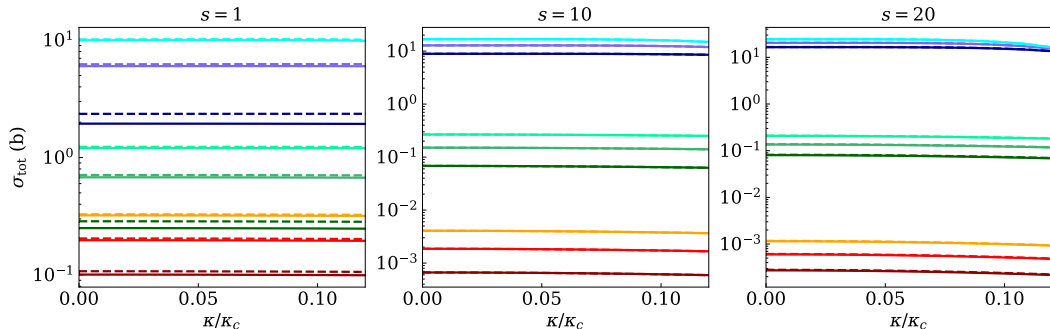
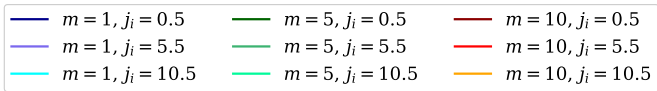
$$\left| S_{fi}^{(1)} \right|^2 \propto \exp\left(-\frac{(\kappa/\kappa_c)^2}{2}\right) \quad (10)$$

$$\kappa_c = m_e \sqrt{H/H_c} \quad (11)$$

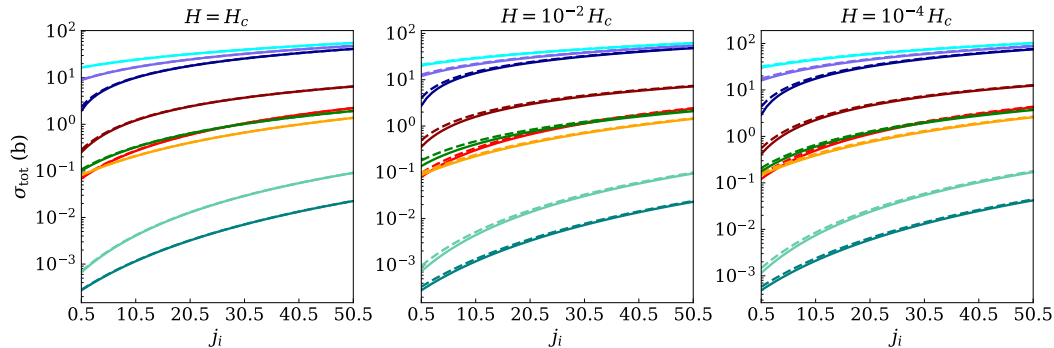
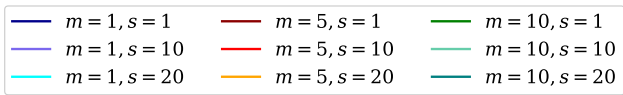
## Luminosity

$$\mathcal{L} = \int d^4x \frac{d^3p}{(2\pi)^3} \frac{d^3k}{(2\pi)^3} v(\mathbf{p}, \mathbf{k}) w^{(e)}(\mathbf{r}, \mathbf{p}) w^{(\gamma)}(\mathbf{r}, \mathbf{k}), \quad (12)$$

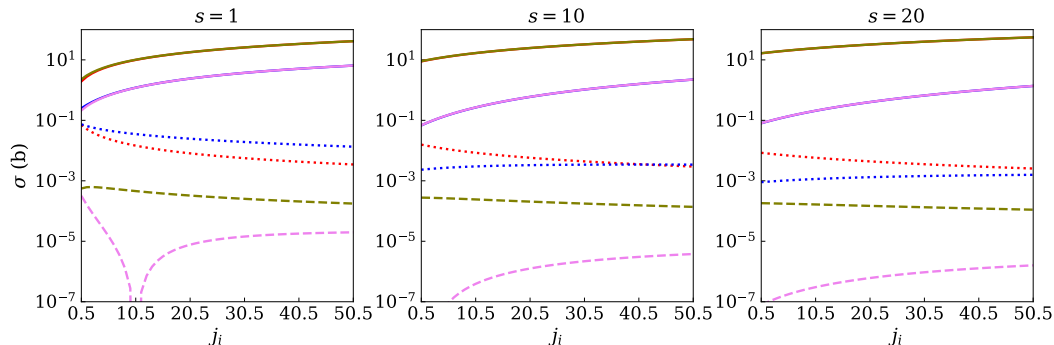
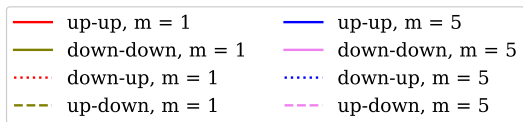
# Total absorption cross-section, $\kappa$



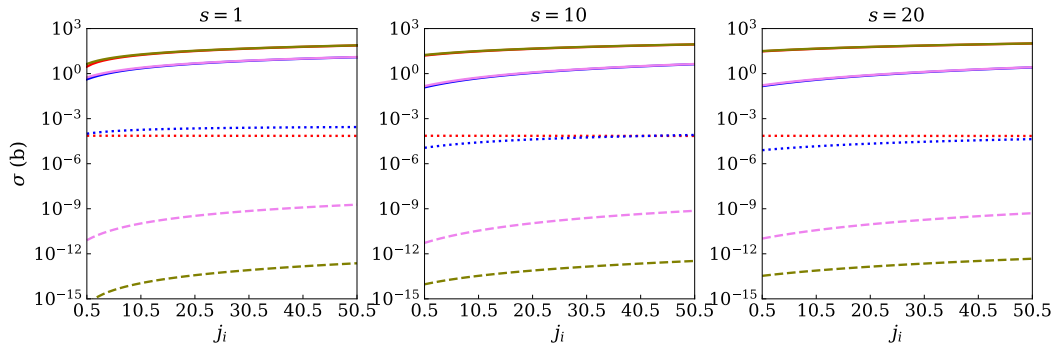
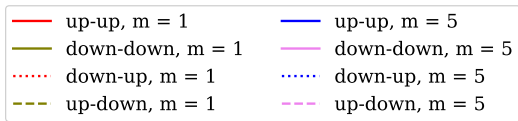
# Total absorption cross-section, $j_i$



# Cross sections for spin transitions, $H = H_c$



# Cross sections for spin transitions, $H = 10^{-4} H_c$



# Key results

Twisted photon absorption cross sections are exponentially suppressed when transverse photon momentum  $\kappa > m_e \sqrt{H/H_c}$ .

Absorption cross sections **grow** with **increasing** TAM of the initial electron and **decrease** with **increasing** TAM of the incident photon

Cross sections for “**down** → **up**” spin flip are always **larger** than for “**up** → **down**” flip. It can be the reminiscent of the Sokolov-Ternov effect.

# Acknowledgments

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# Thank you for your attention!

Reference to our publication:

https:

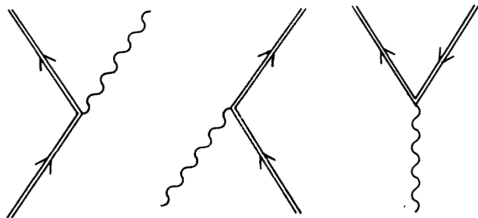
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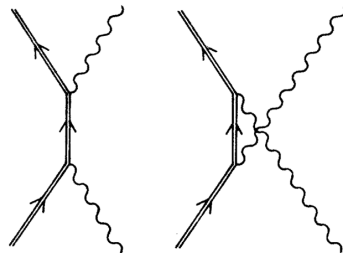


# QED processes in magnetic field

First order



Second order



# Landau states

## Squared Dirac equation

$$\left[ (i\partial_\mu - e\mathcal{A}_\mu)^2 - m_e^2 + eH\Sigma_z \right] \Phi(x) = 0 \quad (13)$$

## Spin operator

$$\Sigma_z = \begin{pmatrix} \sigma_z & 0 \\ 0 & \sigma_z \end{pmatrix} = \text{diag}\{1, -1, 1, -1\}. \quad (14)$$

# Solution to the squared Dirac equation

$$\Phi(x) = \begin{pmatrix} c_1 \Phi_+(x) \\ c_2 \Phi_-(x) \\ c_3 \Phi_+(x) \\ c_4 \Phi_-(x) \end{pmatrix} \quad (15)$$

$$\Phi_{\pm}(x) = N_{s,j} \psi_{s,j \mp 1/2}(\tilde{\rho}) e^{-it\varepsilon_{s,j} + i(j \mp 1/2)\varphi + ip_z z} \quad (16)$$

$$\psi_{s,j \mp 1/2}(\tilde{\rho}) = \tilde{\rho}^{|j \mp 1/2|} L_s^{|j \mp 1/2|}(\tilde{\rho}^2) e^{-\tilde{\rho}^2/2} \quad (17)$$

$$s = 0, 1, 2, \dots$$

$$j = \pm 1/2, \pm 3/2, \dots$$

$$\tilde{\rho} = \rho / \rho_H$$

$$\rho_H = \sqrt{2 / |e| H}.$$

## Solution basis

$$\Phi_{\uparrow}(x) = \begin{pmatrix} \Phi_{+}(x) \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad \sigma_z = \frac{1}{2}, \quad \Phi_{\downarrow}(x) = \begin{pmatrix} 0 \\ \Phi_{-}(x) \\ 0 \\ 0 \end{pmatrix}, \quad \sigma_z = -\frac{1}{2} \quad (18)$$

## Bispinor solution

$$\Psi(x) = [\gamma^{\mu}(i\partial_{\mu} - e\mathcal{A}_{\mu}) + m_e] \Phi(x). \quad (19)$$

# Solution to the Dirac equation

$$\Psi_{\uparrow}(x) = N_{s,j} \begin{pmatrix} (\varepsilon + m_e)\psi_{s,j-1/2}(\tilde{\rho})e^{-i\varphi/2} \\ 0 \\ p_z\psi_{s,j-1/2}(\tilde{\rho})e^{-i\varphi/2} \\ 2\rho_H^{-1}i\psi_{s,j+1/2}(\tilde{\rho})e^{i\varphi/2} \end{pmatrix} e^{i(p_z z + j\varphi - \varepsilon t)}, \quad (20)$$

$$\Psi_{\downarrow}(x) = N_{s,j} \begin{pmatrix} 0 \\ (\varepsilon + m_e)\psi_{s,j+1/2}(\tilde{\rho})e^{i\varphi/2} \\ -2\rho_H^{-1}i(s + j + 1/2)\psi_{s,j-1/2}(\tilde{\rho})e^{-i\varphi/2} \\ -p_z\psi_{s,j+1/2}(\tilde{\rho})e^{i\varphi/2} \end{pmatrix} e^{i(p_z z + j\varphi - \varepsilon t)}. \quad (21)$$

# Vector potential for twisted photon

## Vector potential

$$\mathbf{A}_{\kappa m k_z \Lambda}(\mathbf{r}) = \int a_{\kappa m k_z}(\mathbf{q}) \mathbf{e}_{\mathbf{q}\Lambda} e^{i\mathbf{q}\cdot\mathbf{r} - i\omega t} \frac{d^3 q}{(2\pi)^3}, \quad (22)$$

$$a_{\kappa m k_z}(\mathbf{q}) = \frac{(2\pi)^2}{\kappa} \delta(q_{\perp} - \kappa) \delta(q_z - k_z) e^{im\varphi_q}. \quad (23)$$

## Normalization

$$\int_V \frac{\mathbf{E}^2 + \mathbf{H}^2}{8\pi} d^3 \mathbf{r} = \omega \quad \Longleftrightarrow \quad \int_V |\mathbf{A}|^2 d^3 \mathbf{r} = \frac{2\pi}{\omega} \quad (24)$$

# Vector potential of the twisted photon

$$\mathbf{A}_{\kappa m k_z \Lambda}(\mathbf{r}) = \sqrt{\frac{\pi \kappa}{3\omega LR}} e^{ik_z z - i\omega t} \sum_{\sigma=0, \pm 1} d_{\sigma, \lambda}^{(1)}(\theta_k) e^{i(m-\sigma)\varphi_r} J_{m-\sigma}(\kappa\rho) \boldsymbol{\chi}_\sigma \quad (25)$$

## Wigner matrices

$$d_{\sigma, \lambda}^{(1)}(\theta_k) = \frac{1}{2}(1 + \sigma \lambda \cos \theta_k), \quad \sigma \neq 0$$

$$d_{0, \lambda}^{(1)}(\theta_k) = \frac{\lambda}{2} \sin \theta_k$$

## Spin operator eigenvectors

$$\boldsymbol{\chi}_\pm = \mp \frac{1}{\sqrt{2}} \{1, \pm i, 0\}$$

$$\boldsymbol{\chi}_0 = \{0, 0, 1\}$$

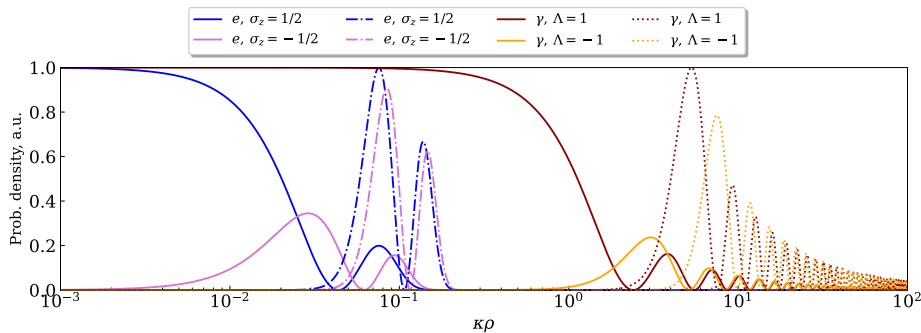
## Vector potential in the coordinate space

$$\mathbf{A}_{\kappa m k_z \Lambda}(\mathbf{r}) = \sqrt{\frac{\pi \kappa}{3\omega LR}} e^{ik_z z + im\varphi_r - i\omega t} \times$$
$$\times \begin{pmatrix} id_{-1\Lambda}^{(1)}(\theta_k) J_{m+1}(\kappa\rho) e^{i\varphi_r} + id_{1\Lambda}^{(1)}(\theta_k) J_{m-1}(\kappa\rho) e^{-i\varphi_r} \\ d_{-1\Lambda}^{(1)}(\theta_k) J_{m+1}(\kappa\rho) e^{i\varphi_r} - d_{1\Lambda}^{(1)}(\theta_k) J_{m-1}(\kappa\rho) e^{-i\varphi_r} \\ \sqrt{2} d_{0\Lambda}^{(1)}(\theta_k) J_m(\kappa\rho) \end{pmatrix}. \quad (26)$$



# Photon and electron density profiles

$s_i = s_f = 1$ ,  $H = H_c$ . Solid lines -  $j = 1/2$ ,  $m = 1$ ; dashed lines -  $j = 11/2$ ,  $m = 5$



# S-matrix for “up-up” transition

$$\begin{aligned}
 S_{\uparrow\uparrow}^{(1)} = & \sqrt{\frac{2\pi\kappa}{3\omega LR}} ie N_i N_f (2\pi)^3 \rho_H \delta(\varepsilon_i - \varepsilon_f + \omega) \delta(p_{zi} - p_{zf} + k_z) \delta_{j_f, j_i+m} \left[ 2\sqrt{2} \times \right. \\
 & \times \left( (\varepsilon_i + m_e) d_{-1\Lambda}^{(1)}(\theta_k) \mathcal{F}_{s_f, s_i}^{j_f+1/2, j_i-1/2}(\kappa\rho_H) - (\varepsilon_f + m_e) d_{1\Lambda}^{(1)}(\theta_k) \mathcal{F}_{s_f, s_i}^{j_f-1/2, j_i+1/2}(\kappa\rho_H) \right) \\
 & \left. + \rho_H d_{0\Lambda}^{(1)}(\theta_k) (m_e(p_{zi} + p_{zf}) + p_{zf}\varepsilon_i + p_{zi}\varepsilon_f) \mathcal{F}_{s_f, s_i}^{j_f-1/2, j_i-1/2}(\kappa\rho_H) \right], \quad (27)
 \end{aligned}$$

$$\mathcal{F}_{s_f, s_i}^{l_f, l_i}(\kappa\rho_H) \propto \exp\left(-(\kappa/\kappa_c)^2/2\right), \quad \kappa_c = \sqrt{2}/\rho_H = m_e \sqrt{H/H_c}$$

## Special function $\mathcal{F}$

$$\mathcal{F}_{s,s'}^{\ell,\ell'}(y) = \int_0^\infty dx x^{\ell+\ell'+1} L_s^\ell(x^2) L_{s'}^{\ell'}(x^2) J_{\ell-\ell'}(yx) e^{-x^2} =$$
$$\frac{(s' + \ell')!}{s!} \frac{1}{2^{2(s-s')+\ell-\ell'+1}} y^{2(s-s')+\ell-\ell'} L_{s'+\ell'}^{s-s'+\ell-\ell'}(y^2/4) L_{s'}^{s-s'}(y^2/4) e^{-y^2/4}, \quad (28)$$

**(Gradshtein & Ryzhik'63)**

Next values are investigated

$$\sigma_{\text{tot}} = \frac{1}{2} \sum_{\Lambda=\pm 1} \sum_{\sigma_{zf}=\pm 1/2} \int d\sigma, \quad (29)$$

$$\sigma_{\text{spin}}^{fi} = \frac{1}{2} \sum_{\Lambda=\pm 1} \int d\sigma. \quad (30)$$

# Luminosity

$$\mathcal{L} = \int d^4x \frac{d^3p}{(2\pi)^3} \frac{d^3k}{(2\pi)^3} v(\mathbf{p}, \mathbf{k}) w^{(e)}(\mathbf{r}, \mathbf{p}) w^{(\gamma)}(\mathbf{r}, \mathbf{k}) \quad (31)$$

$$v(\mathbf{p}, \mathbf{k}) = \frac{p_\mu k^\mu}{\varepsilon(\mathbf{p})\omega(\mathbf{k})} = 1 - \frac{\mathbf{p} \cdot \mathbf{k}}{\varepsilon(\mathbf{p})\omega(\mathbf{k})} \quad (32)$$

# Wigner functions

$$w_{s,j}^{(e)}(\mathbf{r}, \mathbf{p}, t) = \frac{1}{(2\pi)^3} \int \Psi_{s,j}^\dagger(\mathbf{r} - \mathbf{y}/2, t) \Psi_{s,j}(\mathbf{r} + \mathbf{y}/2, t) e^{i\mathbf{p}\cdot\mathbf{r}} d^3\mathbf{y}, \quad (33)$$

$$w^{(\gamma)}(\mathbf{r}, \mathbf{p}, t) = \frac{1}{(2\pi)^3 \omega} \int \mathbf{E}_{\kappa m k_z \Lambda}^*(\mathbf{r} - \mathbf{y}/2, t) \mathbf{E}_{\kappa m k_z \Lambda}(\mathbf{r} + \mathbf{y}/2, t) e^{i\mathbf{k}\cdot\mathbf{r}} d^3\mathbf{y} \quad (34)$$

$$\int w_{s,j}^{(e)}(\mathbf{r}, \mathbf{p}, t) \frac{d^3\mathbf{p}}{(2\pi)^3} = \Psi_{s,j}^\dagger \Psi_{s,j} \quad (35)$$

$$\int w^{(\gamma)}(\mathbf{r}, \mathbf{p}, t) \frac{d^3\mathbf{k}}{(2\pi)^3} = \frac{1}{\omega} \mathbf{E}_{\kappa m k_z \Lambda}^* \mathbf{E}_{\kappa m k_z \Lambda} = \omega \mathbf{A}_{\kappa m k_z \Lambda}^* \mathbf{A}_{\kappa m k_z \Lambda} \quad (36)$$

# Sharp minimum of $\sigma_{\downarrow\uparrow}$ and $\sigma_{\uparrow\downarrow}$

$$\omega_{\downarrow\uparrow,\min} \approx m_e \frac{H}{H_c} \frac{\Lambda}{\Lambda - 1} \frac{2(m+1)}{(1 + \varepsilon/m_e)}, \quad \omega_{\uparrow\downarrow,\min} \approx \frac{\kappa^2}{m_e} \frac{\Lambda}{\Lambda + 1} \frac{(s_i + j_i + 1/2)}{m(1 + \varepsilon_i/m_e)} \quad (37)$$

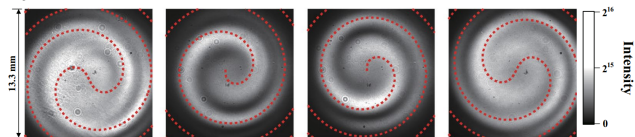
$$j_{\min} \approx \frac{1}{2} \frac{H_c}{H} (m^2 - 1) - s - \frac{1}{2} \quad (38)$$

# Generation of vortices

Electron can emit a vortex photon being

- in an external magnetic field  
(Katoh+'23)
- in a helical undulator  
(Katoh+'23)
- in a circularly polarised laser wave  
(Epp+'23, Zhang+'21)

by helical undulator:



by circularly-polarised laserwave:

