

# Origin of the most energetic particles in the Universe

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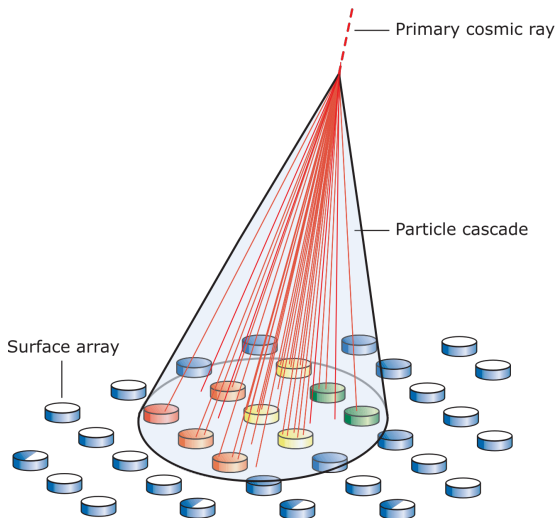
Particles and Cosmology – 2025

при поддержке гранта РФФИ № 24-12-00457

# Lecture 1

- 1 Review of observational data
- 2 Cosmic ray propagation
- 3 Common acceleration mechanisms

# Extensive air shower



Cosmic Rays discovered  
by Victor Hess in 1912

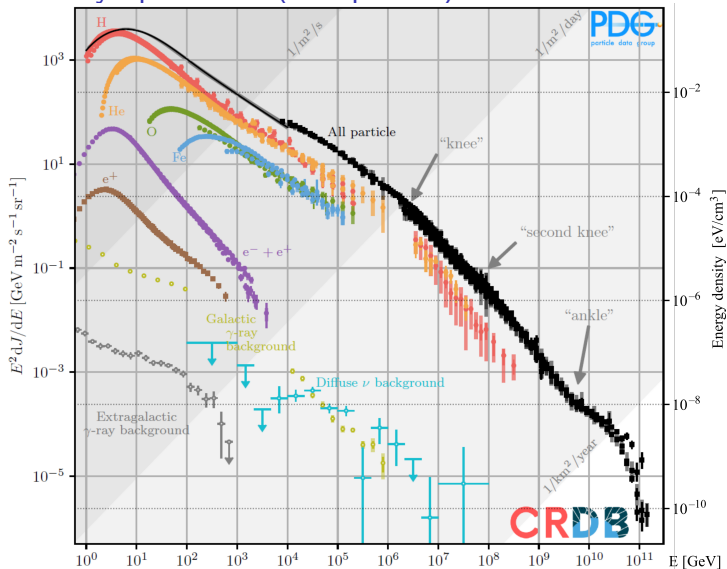
Primary particle  
hits atmosphere  
  
lots of secondaries  
detected on ground

from Bauleo & Martino, Nature 2009

# Pierre Auger Observatory

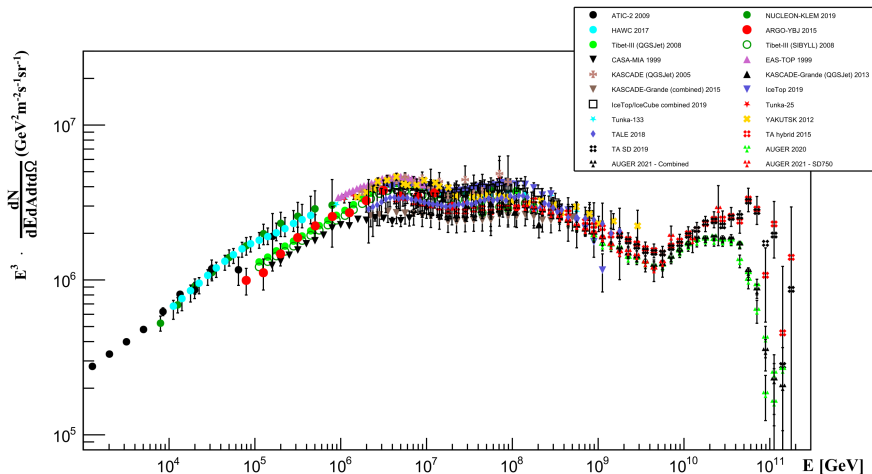


# Cosmic Ray spectrum (all species)



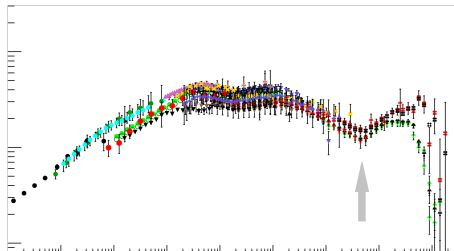
S. Navas et al. (Particle Data Group), Phys. Rev. D 110, 030001 (2024)

# Cosmic Ray spectrum (baryons and nuclei)



Di Sciascio, Appl. Sci. 2022

# Two components in the Cosmic Ray distribution



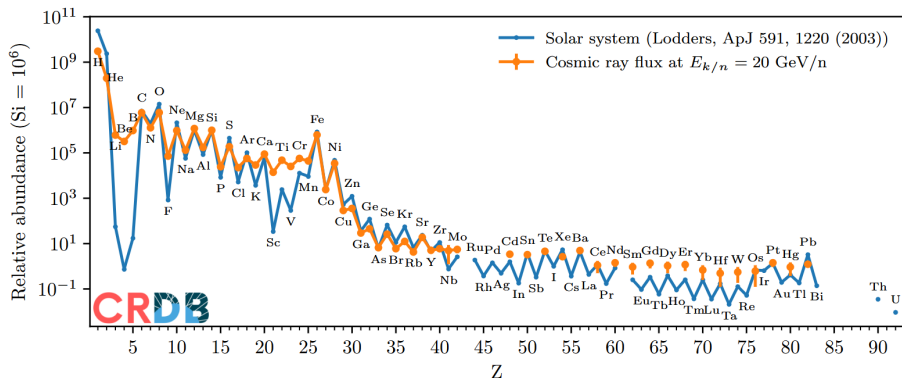
## Lower-energy (likely galactic)

- $E < 5 \times 10^{18} \text{ eV}$
- energy density  $\approx 0.3 \text{ eV/cm}^3$

## Higher-energy (likely extragalactic)

- $E > 5 \times 10^{18} \text{ eV}$
- energy density  $\approx 10^{-8} \text{ eV/cm}^3$

# Cosmic Ray abundances



D. Maurin et al., The European Physical Journal C (2023)

Lithium, Beryllium and Boron nuclei in Cosmic Rays  
are almost entirely of secondary origin



# Nuclei spallation and CR isotopic clocks



$^9\text{Be}$  is stable

$^{10}\text{Be}$  half-life is  $T_{1/2} = 1.39 \times 10^6$  yr (decays into  $^{10}\text{B}$ )

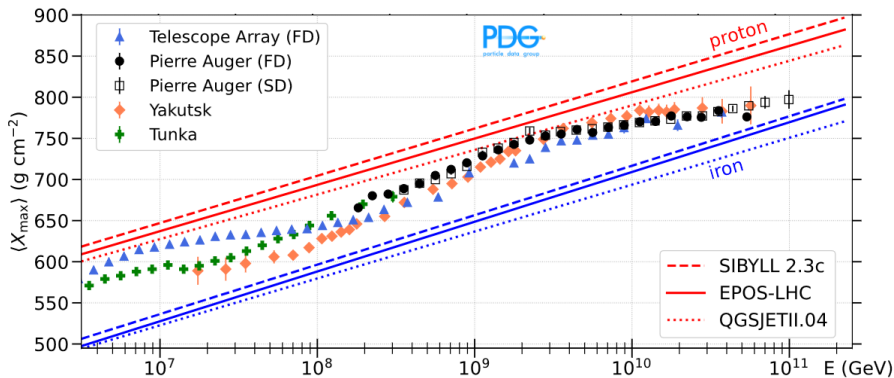
$\text{Be}/\text{O}$  (or  $\text{Be}/\text{C}$ ) ratio measures *grammage*,  $X = \int \rho \, d\ell$

- traversed grammage  $X = 5 \div 7 \text{ g/cm}^2$  (at  $\sim 250 \text{ MeV/nucleon}$ )

$^{10}\text{Be}/^9\text{Be}$  (or  $^{10}\text{Be}/^{10}\text{B}$ ) ratio measures age

- residence time  $t_{\text{res}} \sim 15 \text{ Myr}$  ( $2.5 \div 50 \text{ Myr}$ )

# Cosmic Ray composition at highest energies



S. Navas et al. (Particle Data Group), Phys. Rev. D 110, 030001 (2024)

## Note:

direct measurement of cross-sections is possible only up to  $10^8$  GeV (LHC)

# Confinement by Galactic magnetic field

Lower-energy Cosmic Rays diffuse out of the disc

The diffusion is governed by small-scale turbulent magnetic field

Injection spectrum at the galactic sources  $N_E \propto E^{-\gamma}$

transforms into observed spectrum  $N_E \propto E^{-\gamma-\delta}$  ( $\gamma + \delta \approx 2.7$ )

if escape time scales as  $t_{esc} \propto E^{-\delta}$

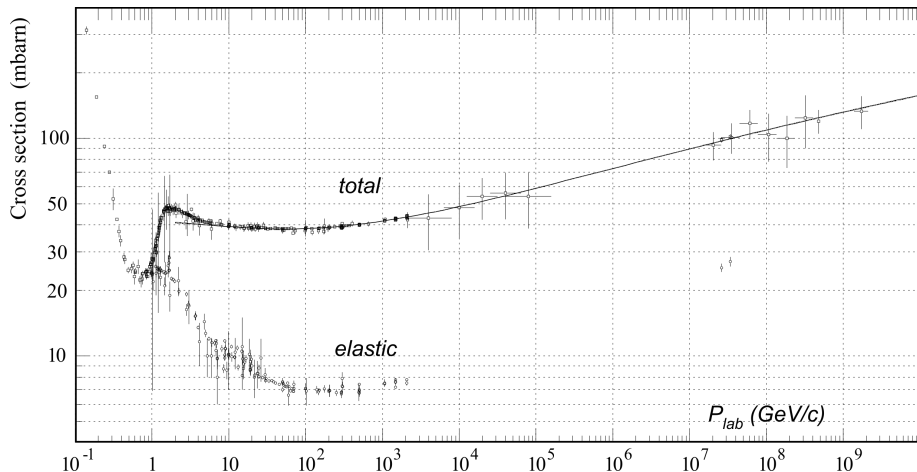
- Bohm-type  $\delta = 1$
- Kolmogorov-type  $\delta = 1/3$
- Kreichnan-type  $\delta = 1/2$

# Deflection by Galactic magnetic field

- large-scale magnetic field strength in Galactic disk  $\sim 3\mu\text{G}$
- thickness of Galactic disk  $\sim 300$  parsec
- above  $\sim 2 \times 10^{19}$  eV all Cosmic Rays are not confined in the disc
- extragalactic Cosmic Rays deflect by an angle  $\theta_{\text{defl}} \sim Z/E_{18}$

Arrival directions of the highest-energy CRs ( $\gtrsim 10^{20}$  eV)  
point towards sources with less than  $15^\circ$  deviation even for iron nuclei  
(less than  $1^\circ$  for protons)

# proton-proton cross-section



Olive et al. (Particle Data Group), Chin. Phys. C (2014)

# Interaction of Cosmic Rays with matter

## pp reaction

$p + p \rightarrow 2 \text{ nucleons} + \text{pions}$  (main channel)

cross-section  $\approx 30 \text{ mbarn}$ , almost independent of energy in GeV-TeV range

Three branches with approx. equal probabilities,  $\pi^0, \pi^+, \pi^-$

$$\pi^0 \rightarrow 2\gamma$$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e$$

$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$$

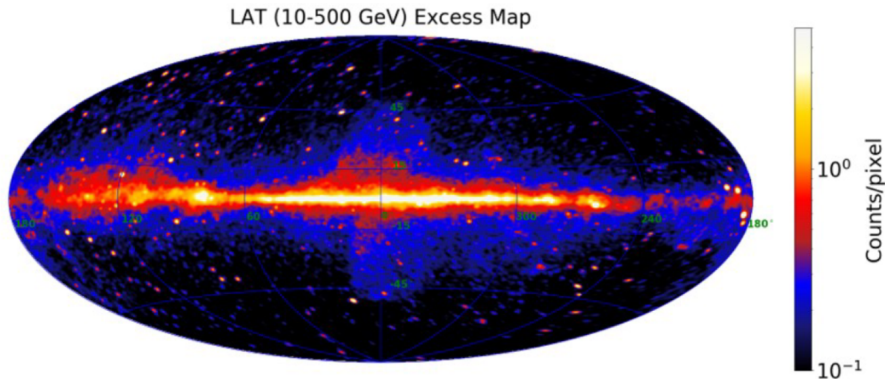
$$\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$$

## energy shares in secondaries

- (1/6): electrons and positrons
- (1/3): photons
- (1/2): neutrinos  $\leftarrow$  source of astrophysical neutrinos (1 of 2)

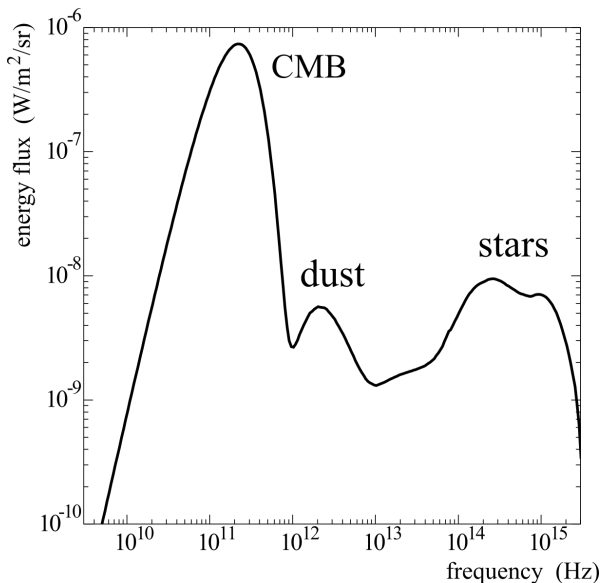
Each of these particles has energy  $\sim 0.1$  of the primary's energy

# GeV sky map



Workman et al. (Particle Data Group), Prog. Theor. Exp. Phys. (2022)

# Extragalactic background light



Stecker & Salamon, ApJ 1999



# Energy loss channels

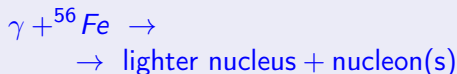
## $e^-e^+$ -pair production (Bethe-Heitler process)



cross-section  $\sigma_{pp} \approx 1.5 Z^2 \text{ mbarn}$

inelasticity  $\epsilon \approx 10^{-4} A^{-1}$

## photodisintegration



cross-section  $\sigma_{\text{dis}} \approx 100 \text{ mbarn}$

inelasticity  $\epsilon \approx 0.02$

## photo-pion reaction

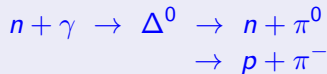


cross-section  $\sigma_{\pi} \approx 0.6 \text{ mbarn}$

inelasticity  $\epsilon \approx 0.2$

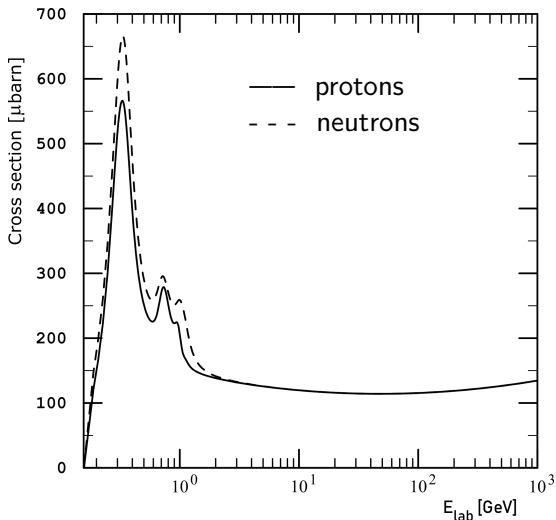
Inelasticity  $\epsilon$  — fraction of energy lost per interaction

# Photo-pion reactions



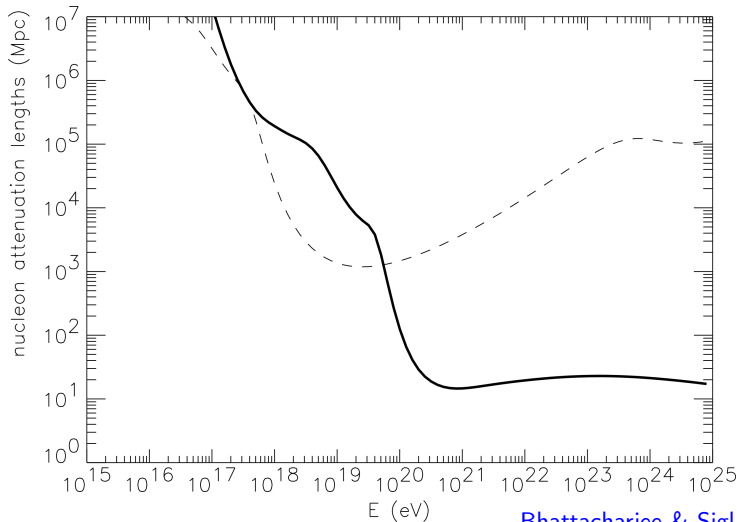
decay of charged pions  
→ 2nd source  
of astrophysical neutrinos

$$E_\nu \sim 0.05 E_p$$



Bhattacharjee & Sigl, Phys. Rep. 2000

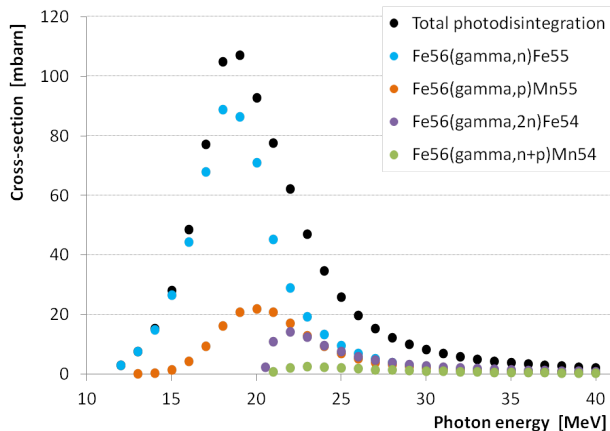
# GZK (Greisen–Zatsepin–Kuzmin) horizon



Bhattacharjee & Sigl, Phys. Rep. 2000

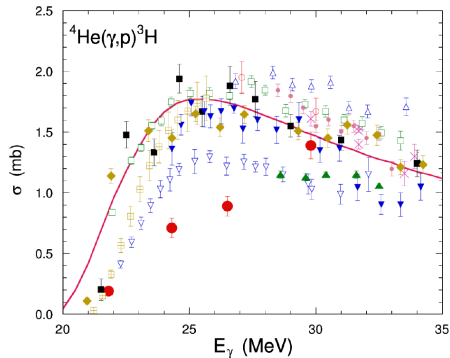
nucleon attenuation lengths: solid line — photo-pion production  
dashed line — pair production

# Giant Dipole Resonance photo-disintegration

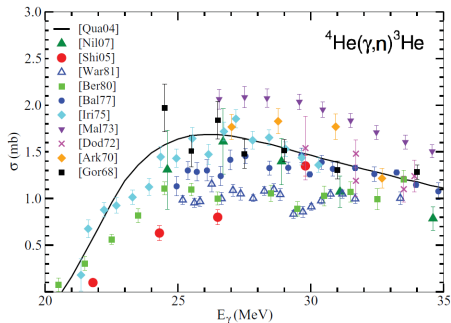


Borodina et al. (2000)

# Photo-disintegration of Helium



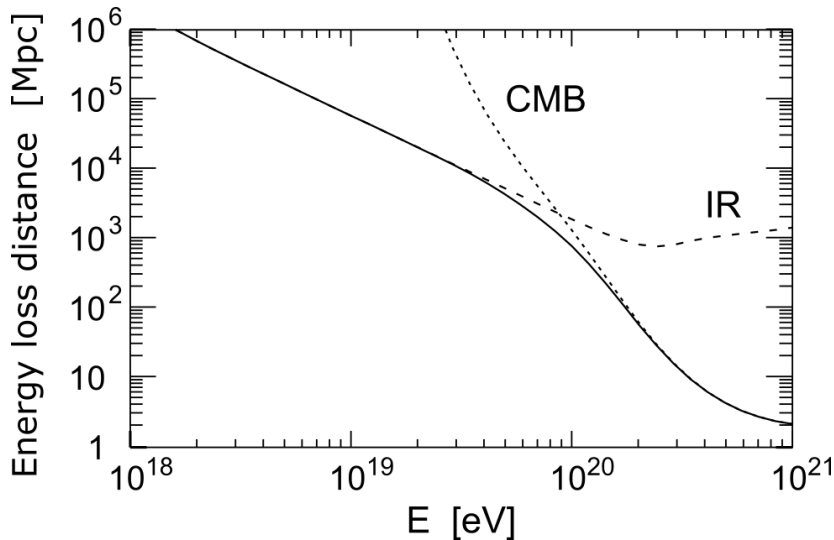
Raut et al., Phys Rev Lett (2012)



Tornow et al., Phys Rev C (2012)

For He nuclei, photodisintegration dominates over pair production losses

## Iron nuclei losses



Epele & Roulet, Phys. Rev. Lett. 1998

# Cosmic Ray horizon at $E = 10^{20}$ eV

## protons

Attenuation distance  $\sim 100$  Mpc (due to photo-pion losses)

- $\sim 30$  AGNs and a few galaxy clusters are within this distance

## Helium nuclei

Attenuation distance  $\sim 2$  Mpc (due to photodisintegration)

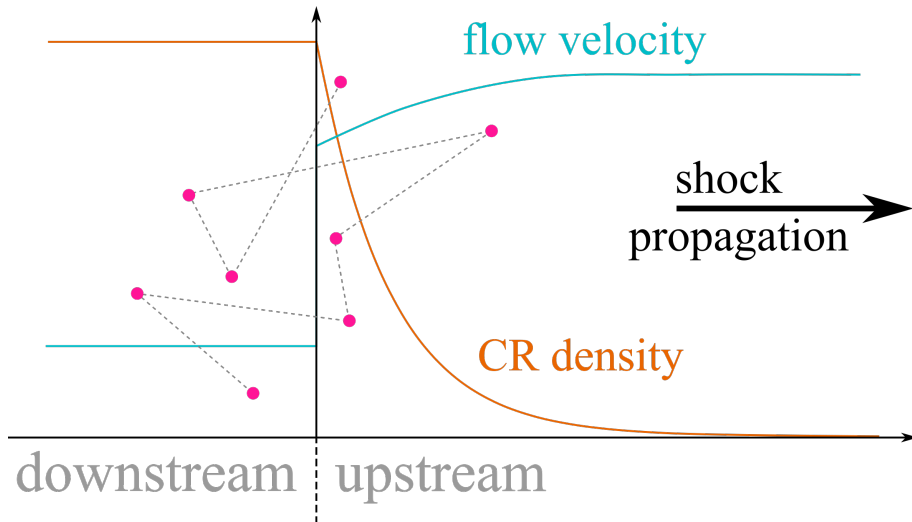
- the nearest *currently active* AGN (Cen A) is at  $3 \div 4$  Mpc
- the nearest galaxy cluster (the Virgo Cluster) is at  $15 \div 20$  Mpc

## Iron nuclei

Attenuation distance  $\sim 700$  Mpc (due to photodisintegration)

- $\sim 10^4$  AGNs and  $\sim 300$  galaxy clusters are within this distance

# Schematic view of particle-accelerating shock





# Diffusive shock acceleration

Flow's work done on CRs

$$P_{flow} = \left( \frac{1}{3} E N_E dE \right) V_{sh}$$

=

Power of acceleration

$$P_{acc} = \dot{E} (N_E dE \lambda)$$

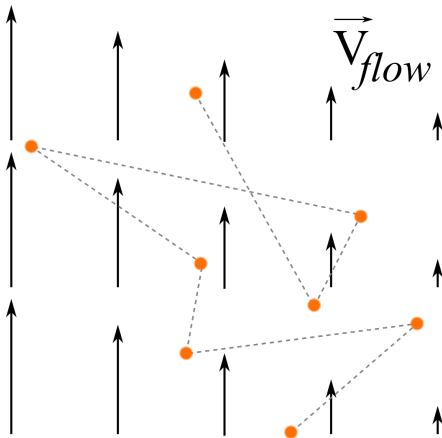
Acceleration rate

$$\dot{E} = \frac{1}{3} \frac{E V_{sh}}{\lambda} = \frac{1}{3} \frac{E V_{sh}^2}{\mathcal{D}} \left( = \frac{V_{sh}^2}{c^2} q B c \text{ with Bohm diffusion} \right)$$

Energy distribution — power-law  $N_E \propto E^{-\alpha}$  with index  $\alpha = \frac{(V_u/V_d) + 2}{(V_u/V_d) - 1}$

In strong non-relativistic shocks compression ratio is  $V_u/V_d = 4$  and  $\alpha = 2$

# Acceleration in shear flows



Acceleration rate

$$\dot{E} \sim \frac{(\nabla V_{fl} \lambda)^2}{c^2} q B c \propto E^2$$

**innate problem:**

acceleration takes too long  
at lower energies ( $t_{acc} \propto E^{-1}$ )

more like an energy-boost process  
rather than  
standalone acceleration mechanism

# Summary 1

- Cosmic Rays extend in energy beyond  $10^{20}$  eV.
- There are two components:
  - lower-energy, likely Galactic, with  $E_{\text{CR}} < 5 \times 10^{18}$  eV
  - higher-energy, likely extragalactic, with  $E_{\text{CR}} > 5 \times 10^{18}$  eV
- The highest-energy cosmic rays ( $E_{\text{CR}} \gtrsim 10^{20}$  eV) must come from nearby sources
- There are two common (diffusive) acceleration mechanisms:
  - shock acceleration
  - shear-flow acceleration acceleration (energy boost)
- *One more mechanism will be discussed later*